

Understanding orbital physics in correlated materials:

insights from
the DFT+DMFT approach

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Peter Grünberg Institut

Forschungszentrum Jülich

what are correlated materials?

systems for which *single-electron picture* fails *qualitatively*
atomic-like physics survives in the solid

H																	He
Li	Be											B	C	N	O	F	Ne
Na	Mg											Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe
Cs	Ba	• Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn
Fr	Ra	••• Lr	Rf	Db	Sg	Bh	Hs	Mt									

• La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb
•• Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No

Coulomb-induced metal-insulator transition
heavy-Fermions
unconventional superconductivity
spin-charge separation

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Na	Mg											Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe
Cs	Ba	● Lu	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn
Fr	Ra	●● Lr	Rf	Db	Sg	Bh	Hs	Mt									
		● La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb		
		●● Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No		

bad news: DFT (Kohn-Sham) band-structure
=single-electron picture

failure of single-electron picture: *more is different*

Born-Oppenheimer approximation, non-relativistic

$$\hat{H}_e = \underbrace{-\frac{1}{2} \sum_i \nabla_i^2}_{\text{i=electrons}} + \underbrace{\frac{1}{2} \sum_{i \neq i'} \frac{1}{|\mathbf{r}_i - \mathbf{r}_{i'}|}}_{\text{electron-electron interactions}} - \sum_{i,\alpha} \frac{Z_\alpha}{|\mathbf{r}_i - \mathbf{R}_\alpha|} + \frac{1}{2} \sum_{\alpha \neq \alpha'} \frac{Z_\alpha Z_{\alpha'}}{|\mathbf{R}_\alpha - \mathbf{R}_{\alpha'}|}$$

$\alpha = \text{nuclei}$

simple interactions among many particles
can lead to unexpected *emergent* co-operative behavior

Coulomb-induced metal-insulator transition
heavy-Fermions
unconventional superconductivity
spin-charge separation

.....

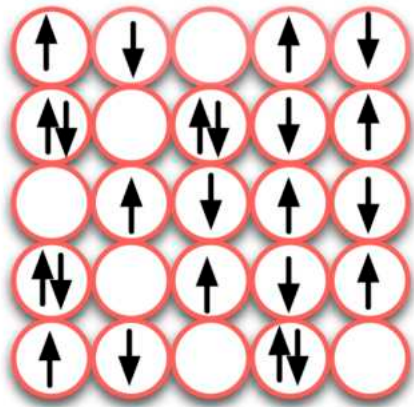


Philip Warren Anderson

4 August 1972, Volume 177, Number 4047

SCIENCE

the key phenomenon: Mott transition



Hubbard model

half filling

$$\hat{H} = \varepsilon_d \sum_i \sum_{\sigma} c_{i\sigma}^{\dagger} c_{i\sigma} - t \sum_{\langle ii' \rangle} \sum_{\sigma} c_{i\sigma}^{\dagger} c_{i'\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

$U/W=0$:
metal

$$m^* \rightarrow \infty$$

metal with heavy
quasiparticles
(large masses)

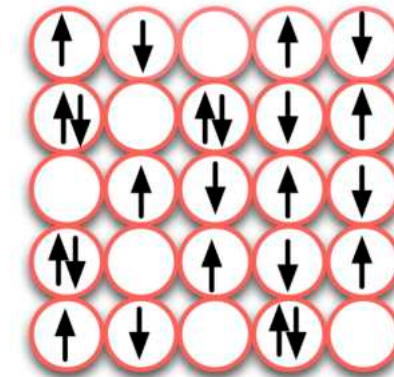
$$\Delta E \rightarrow U - W$$

insulator with
increasing gap

$W/U=0$:
insulator

dynamical mean-field theory (DMFT)

$$\hat{H} = \varepsilon_d \sum_i \sum_{\sigma} c_{i\sigma}^{\dagger} c_{i\sigma} - t \sum_{\langle ii' \rangle} \sum_{\sigma} c_{i\sigma}^{\dagger} c_{i'\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



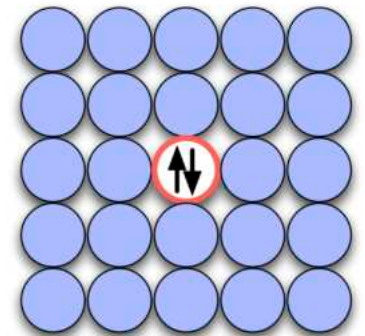
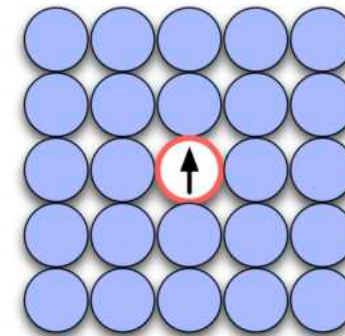
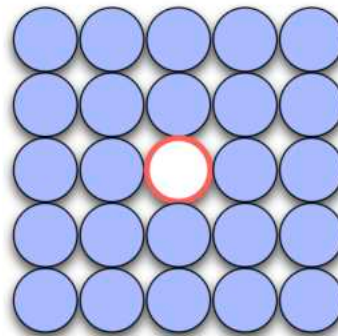
$G^{i,j}$

H^{LDA}

$U^{i,i}$



self-consistent
quantum-impurity model



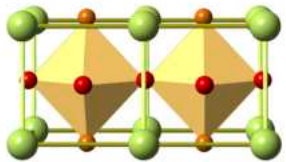
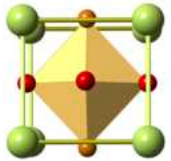
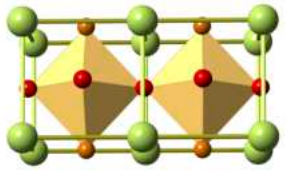
k-independent self-energy

exact in the infinite coordination number limit

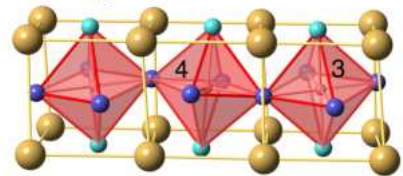
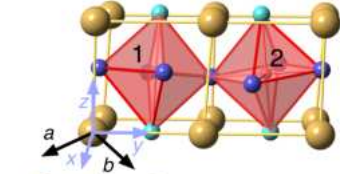
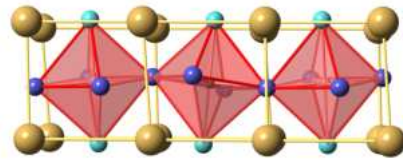
$$\mathcal{G}^{-1} = G^{-1} + \Sigma$$

$$G = G^{i,i}$$

correlated materials: not only W/U

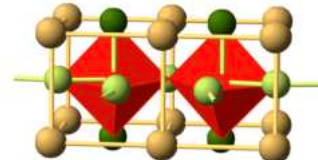
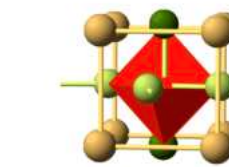
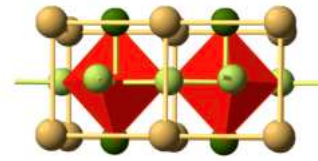


metal



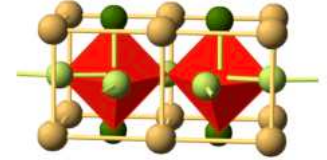
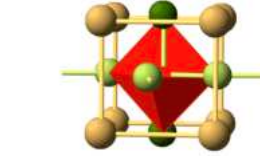
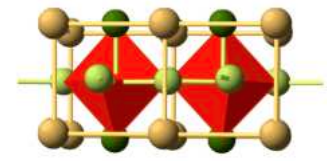
insulator

$(t_{2g})^4$



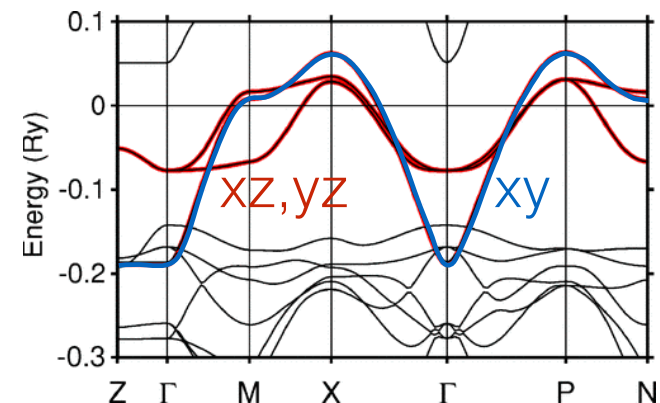
metal

$(t_{2g})^5$



insulator

orbital physics: interplay of spin, orbital, charge, lattice degrees of freedom in the presence of correlations

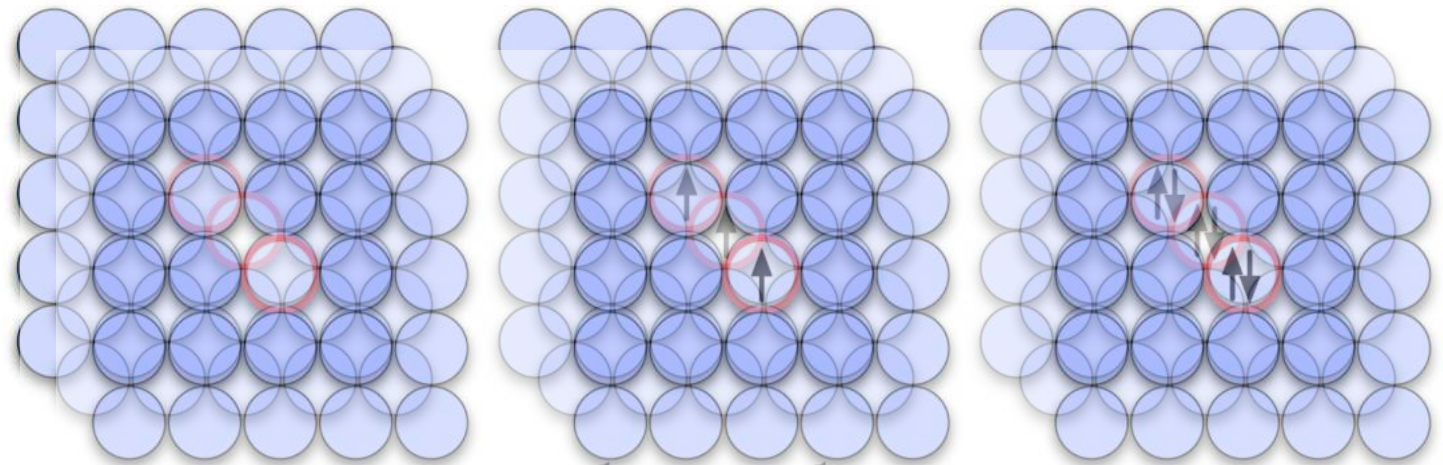
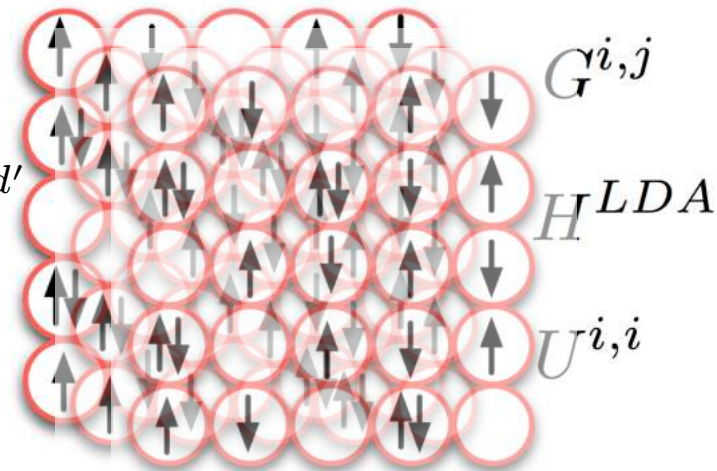


DFT+DMFT: DMFT for Materials

$$\hat{H}_e = \sum_{ab} t_{ab} c_a^\dagger c_b + \frac{1}{2} \sum_{cdc'd'} U_{cdd'c'} c_c^\dagger c_d^\dagger c_{c'} c_{d'}$$



realistic
self-consistent
quantum-impurity
(QI) model



we need **minimal** material-specific models
and **flexible and efficient** solvers

minimal models: DFT-based Wannier functions

VOLUME 92, NUMBER 17

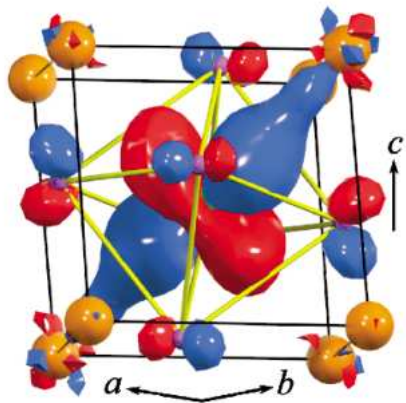
PHYSICAL REVIEW LETTERS

week ending
30 APRIL 2004

Mott Transition and Suppression of Orbital Fluctuations in Orthorhombic $3d^1$ Perovskites

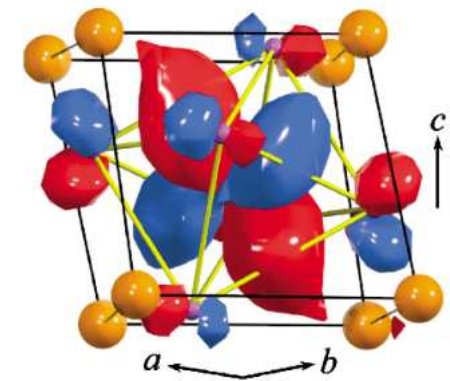
t_{2g}^1

E. Pavarini,¹ S. Biermann,² A. Poteryaev,³ A. I. Lichtenstein,³ A. Georges,² and O. K. Andersen⁴

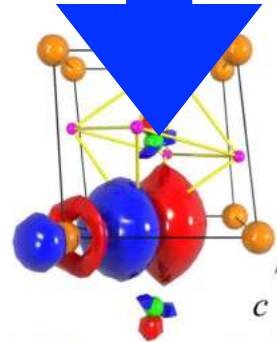


LaTiO_3

DFT+DMFT



YTiO_3



but screened U,J: cRPA,cLDA

flexible and efficient solvers: our package

$$\begin{aligned}
 H = & - \sum_{ii'} \sum_{mm'} \sum_{\sigma} t_{mm'}^{ii'} c_{im\sigma}^{\dagger} c_{i'm'\sigma} \\
 & + U \sum_{im} n_{im\uparrow} n_{im\downarrow} \\
 & + \frac{1}{2} \sum_{im \neq m' \sigma \sigma'} (U - 2J - J\delta_{\sigma\sigma'}) n_{im\sigma} n_{im'\sigma'} \\
 & - J \sum_{m \neq m'} (c_{m\uparrow}^{\dagger} c_{m'\downarrow}^{\dagger} c_{m'\uparrow} c_{m\downarrow} + c_{m\uparrow}^{\dagger} c_{m\downarrow}^{\dagger} c_{m'\uparrow} c_{m'\downarrow})
 \end{aligned}$$

self-energy matrix in spin-orbital space

DMFT and cDMFT

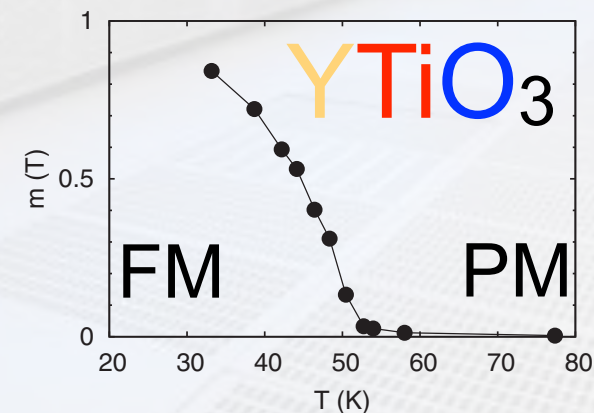
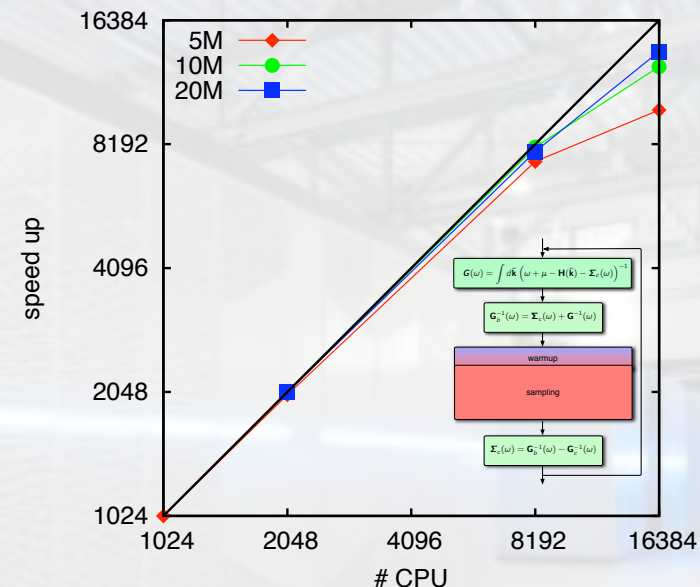
generalized quantum impurity solvers:

general HF QMC

general CT-INT QMC

general CT-HYB QMC

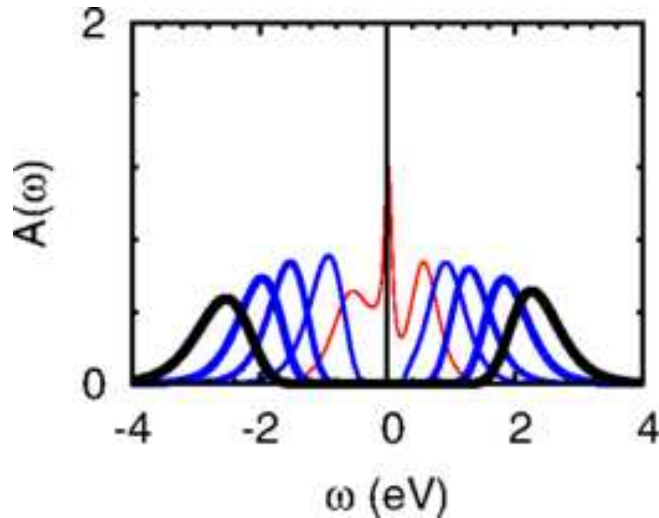
- ♦ CT-HYB: A. Flesch, E. Gorelov, E. Koch and E. Pavarini
[Phys. Rev. B **87**, 195141 \(2013\)](#)
- ♦ CT-INT: E. Gorelov et al, [PRL **104**, 226410 \(2010\)](#)
- ♦ CT-INT+SO: G. Zhang, E. Gorelov, E. Sarvestani, and E. Pavarini,
[Phys. Rev. Lett. **116**, 106402 \(2016\)](#)



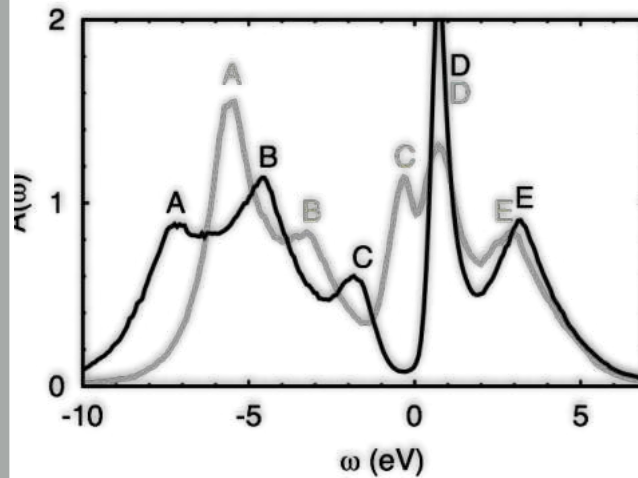
sign problem: smart adapted basis choice

what can we do so far? a lot!

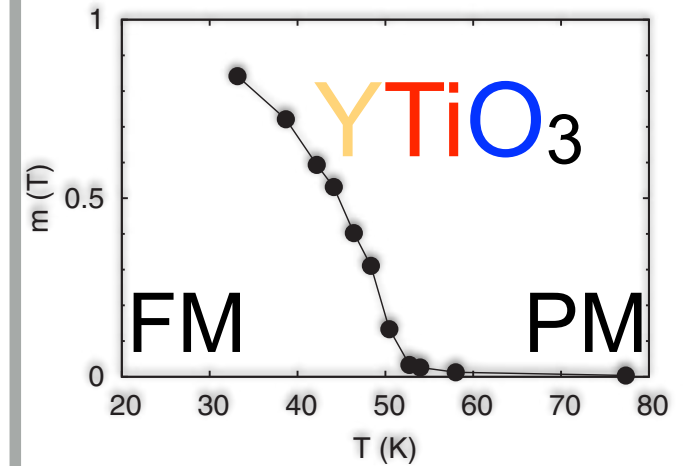
spectral functions



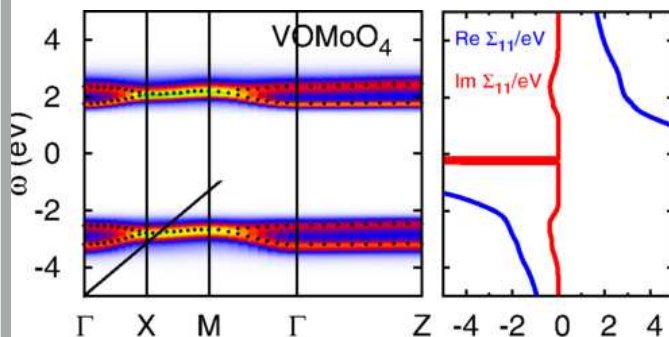
many orbitals



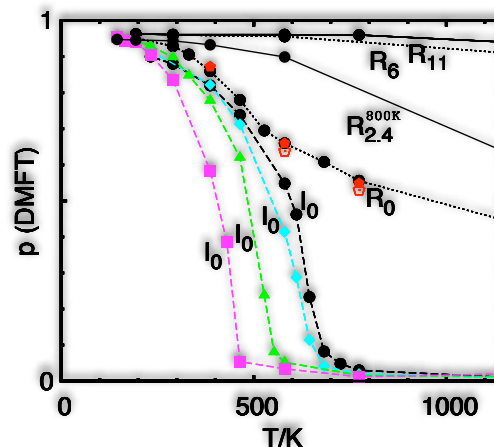
low T



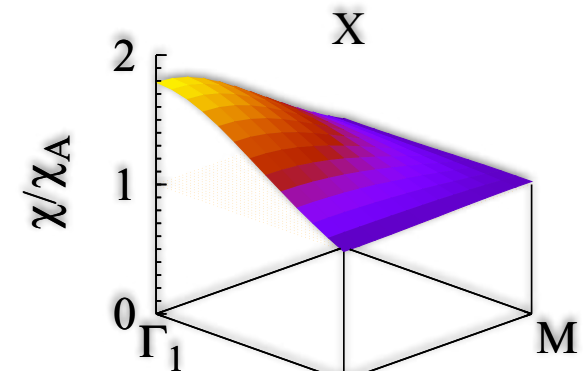
correlated bands



phase transitions

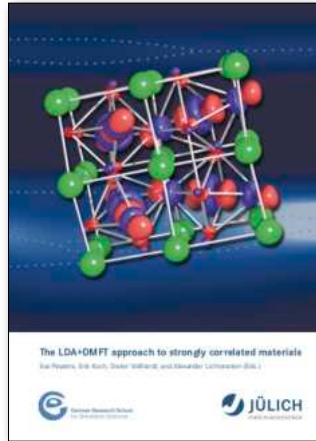


susceptibilities



want to know more?

Autumn School on Correlated Electrons



La Rivista del Nuovo Cimento (2021) 44:597–640
<https://doi.org/10.1007/s40766-021-00025-8>



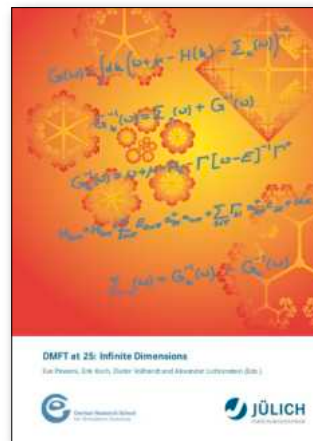
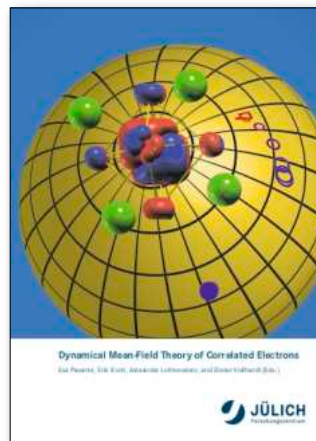
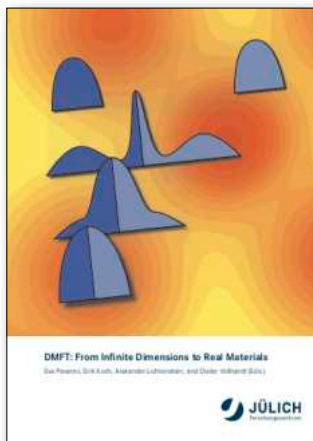
REVIEW PAPER



Solving the strong-correlation problem in materials

Eva Pavarini¹

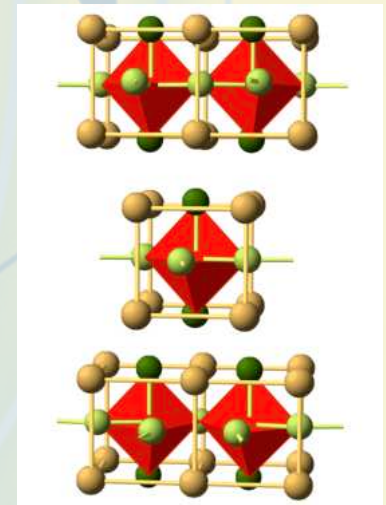
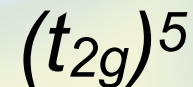
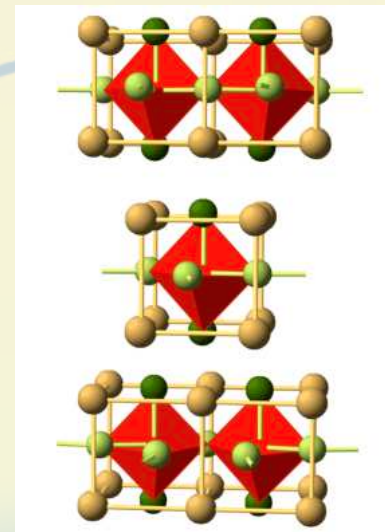
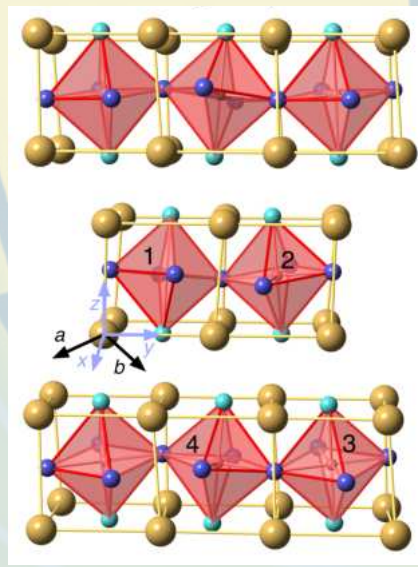
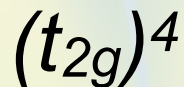
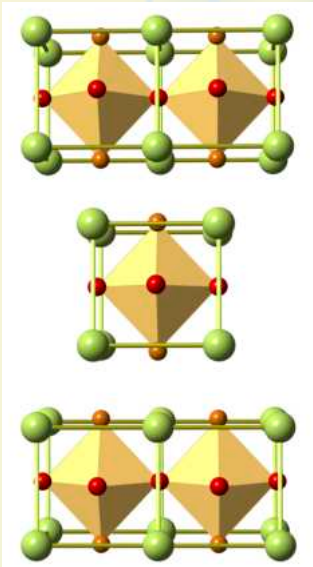
Received: 24 March 2021 / Accepted: 11 June 2021 / Published online: 14 July 2021
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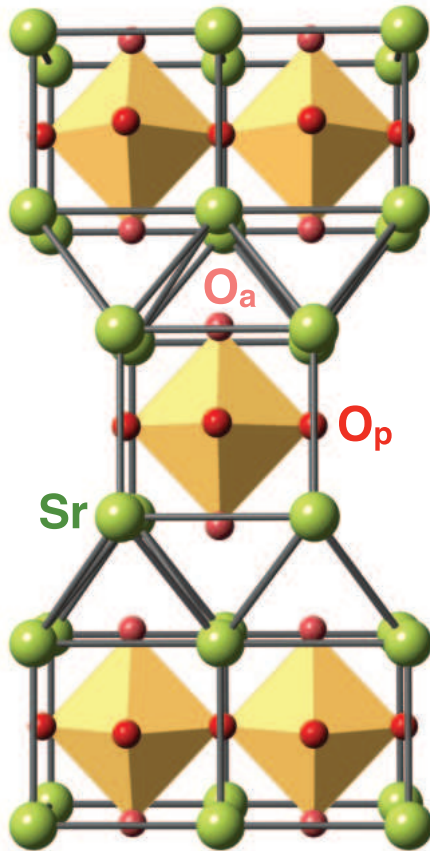
www.cond-mat.de/events/correl.html

Sr₂RuO₄ and its family

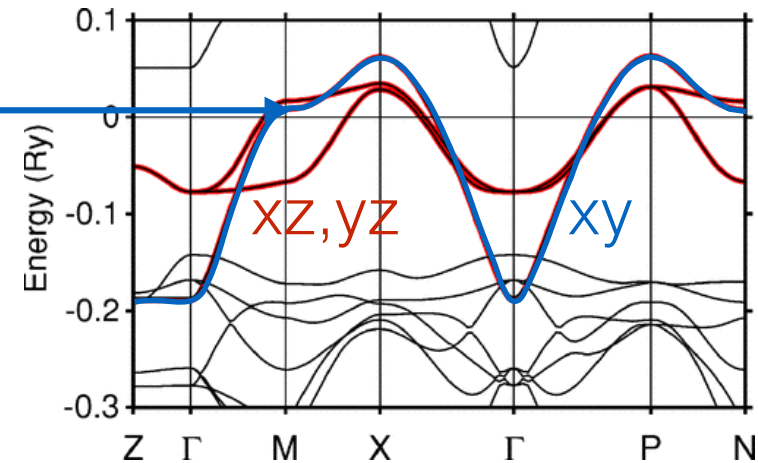
disentangling complexity
building a unifying picture



Sr₂RuO₄: a paradigmatic system



xy vHS



$$H = - \sum_{ii'} \sum_{mm'} \sum_{\sigma\sigma'} t_{m\sigma, m'\sigma'}^{i, i'} c_{im\sigma}^\dagger c_{i'm'\sigma'} + \frac{1}{2} \sum_i \sum_{mm'pp'} \sum_{\sigma\sigma'} U_{mm'pp'} c_{im\sigma}^\dagger c_{im'\sigma'}^\dagger c_{ip'\sigma'} c_{ip\sigma}$$

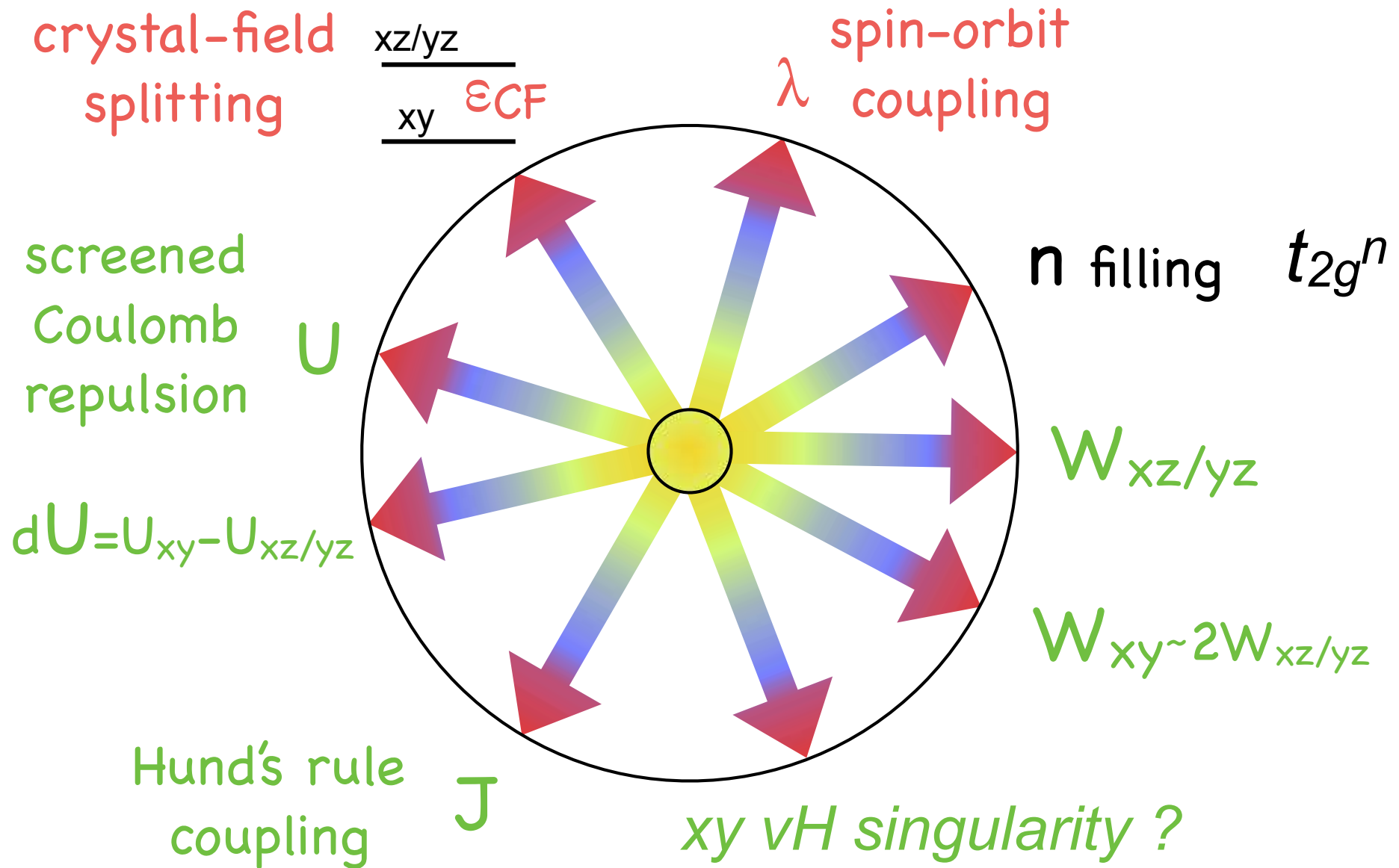
Ru t_{2g}⁴

xy-band 2-dimensional

xz/yz-bands 1-dimensional

$$W_{xy} \sim 2W_{xz/yz}$$

many energy scales are similar



minimal Hubbard model: t_{2g} bands

$$H = - \sum_{ii'} \sum_{mm'} \sum_{\sigma\sigma'} t_{m\sigma, m'\sigma'}^{i, i'} c_{im\sigma}^\dagger c_{i'm'\sigma'} \\ + \frac{1}{2} \sum_i \sum_{mm'} \sum_{pp'} \sum_{\sigma\sigma'} U_{mm'pp'} c_{im\sigma}^\dagger c_{im'\sigma'}^\dagger c_{ip'\sigma'} c_{ip\sigma}$$

$$H_{\text{SO}} = \sum_{i\mu} H_{\text{SO}}^{i\mu} = \sum_{i\mu} \sum_{m\sigma m'\sigma'} \lambda_\mu^i \xi_{m\sigma m'\sigma'}^{i\mu} c_{im\sigma}^\dagger c_{im'\sigma'}$$

local spin-orbit interaction

example: Hund's rule coupling J & van-Hove singularity

PRL **106**, 096401 (2011)

PHYSICAL REVIEW LETTERS

week ending
4 MARCH 2011

Coherence-Incoherence Crossover and the Mass-Renormalization Puzzles in Sr_2RuO_4

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⁶Physics Department and Center for Materials Theory, Rutgers University, Piscataway New Jersey 08854, USA

⁷Collège de France, 11 place Marcelin Berthelot, 75005 Paris, France

(Received 27 October 2010; published 2 March 2011)

We calculate the electronic structure of Sr_2RuO_4 , treating correlations within dynamical mean-field theory. The approach successfully reproduces several experimental results and explains the key properties of this material: the anisotropic mass renormalization of quasiparticles and the crossover into an incoherent regime above a low temperature scale. While the orbital differentiation originates from the proximity of the van Hove singularity, strong correlations are caused by the Hund's coupling. The generality of this mechanism for other correlated materials is pointed out.

Hund's metal

*xy vH singularity
and mass inversions*

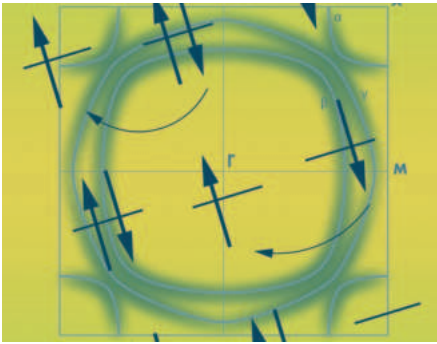
$$W_{xy} > W_{xz}$$

$$\text{but } m^*_{xy} > m^*_{xz}$$

TABLE I. Mass enhancement of the xy and xz orbitals, as a function of Hund's coupling, for $U = 2.3$ eV. Other columns: coherence temperatures as defined in the text.

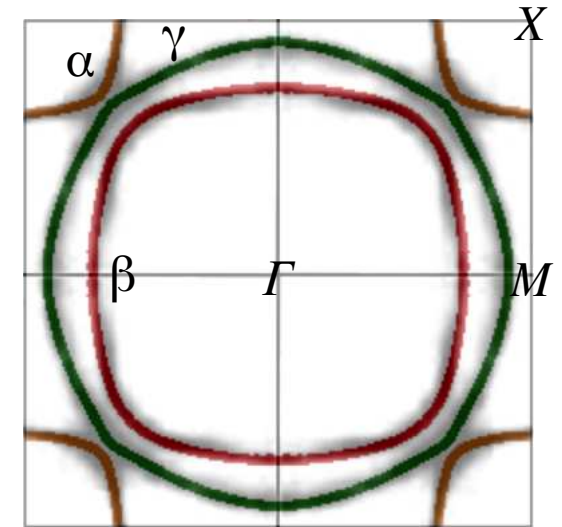
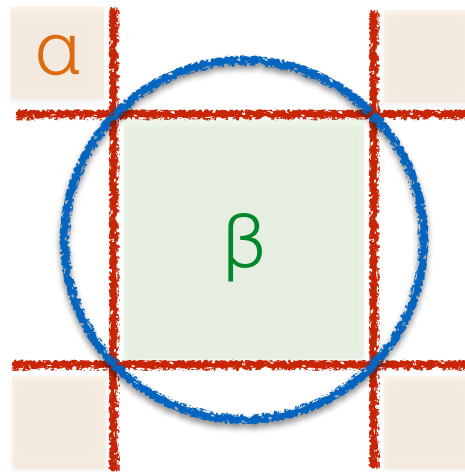
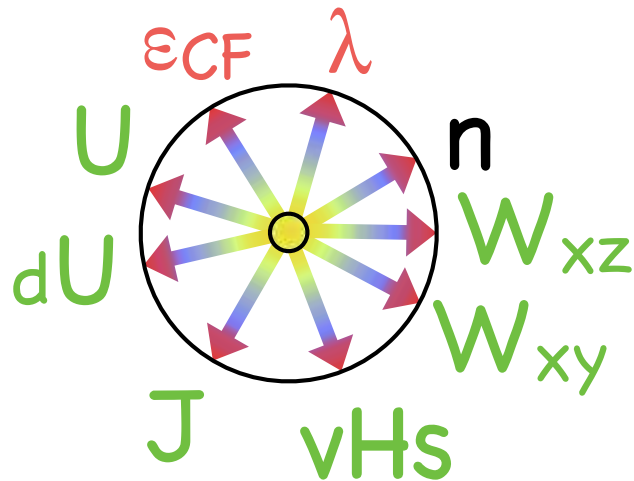
J [eV]	$m^*/m_{\text{LDA}} _{xy}$	$m^*/m_{\text{LDA}} _{xz}$	T^*_{xy} [K]	T^*_{xz} [K]	$T_{>}$ [K]
0.0, 0.1	1.7	1.7	>1000	>1000	>1000
0.2	2.3	2.0	300	800	>1000
0.3	3.2	2.4	100	300	500
0.4	4.5	3.3	60	150	350

example: the Fermi surface puzzle



$$\epsilon_{CF} + \Delta \epsilon_{CF}(U) - \Delta' \epsilon_{CF}(dU) \sim \epsilon_{CF}$$

$$\lambda + \Delta \lambda \sim 2\lambda$$



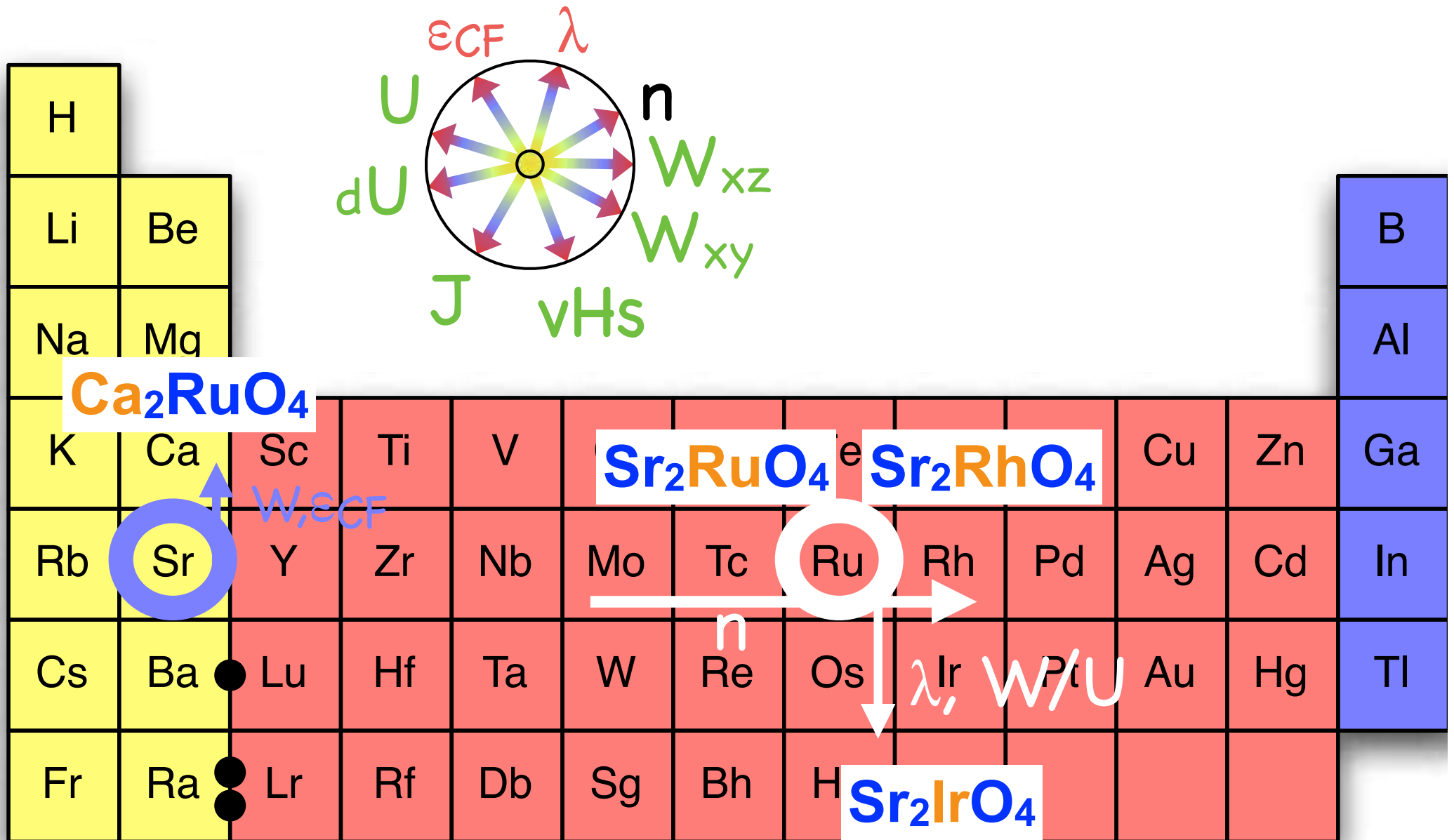
theory:

G. Zhang, E. Gorelov, E. Sarvestani, and E. Pavarini, Phys. Rev. Lett. **116**, 106402 (2016)

experiments:

A. Damascelli, D. H. Lu, K. M. Shen, N. P. Armitage, F. Ronning, D. L. Feng, C. Kim, Z.-X. Shen, T. Kimura, Y. Tokura, Z. Q. Mao, and Y. Maeno, Phys. Rev. Lett. **85**, 5194 (2000)

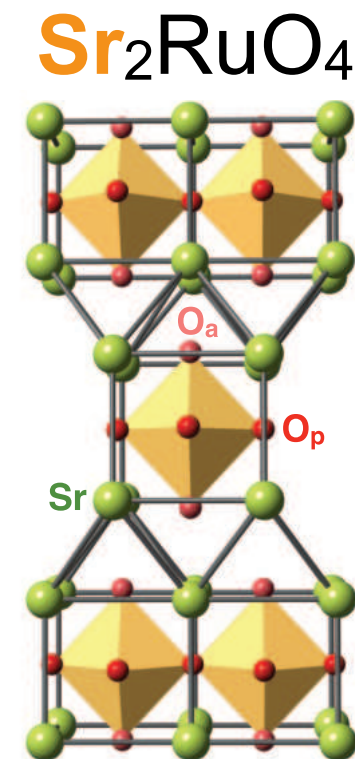
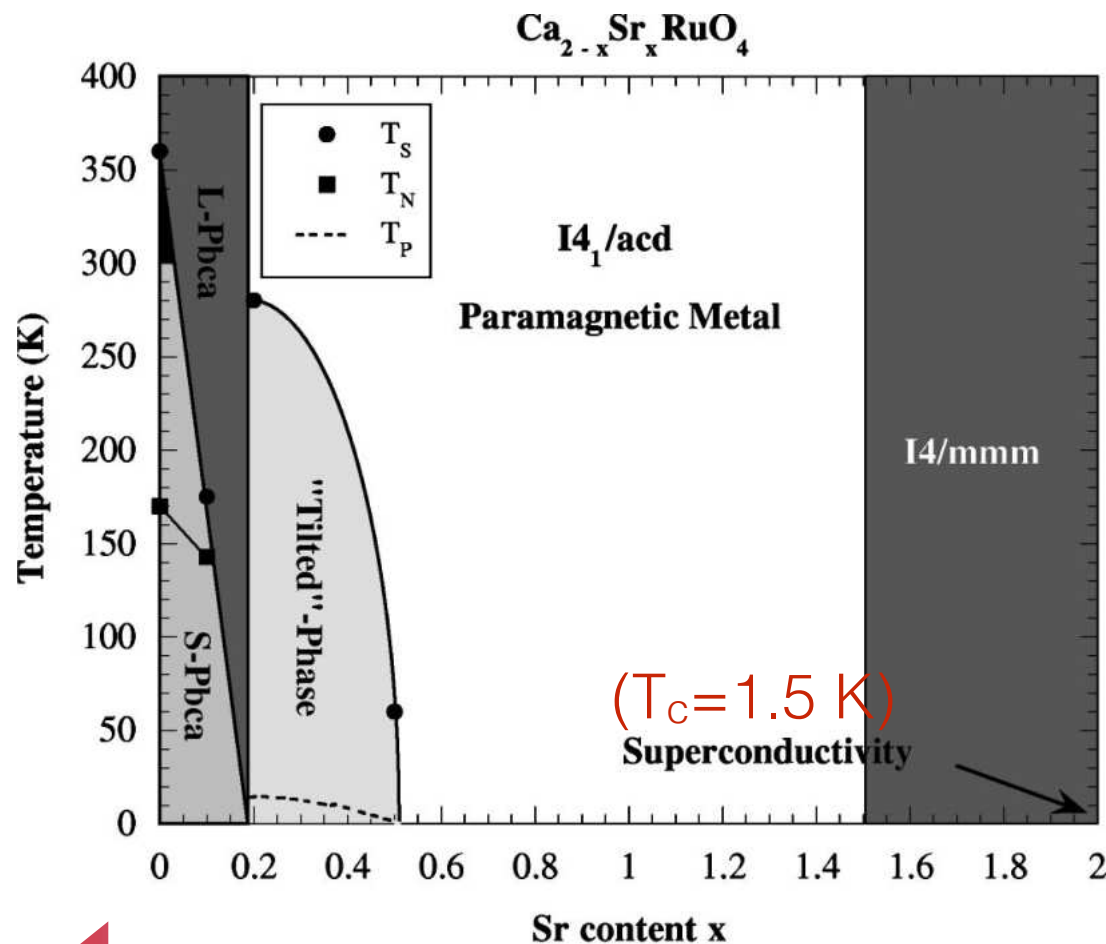
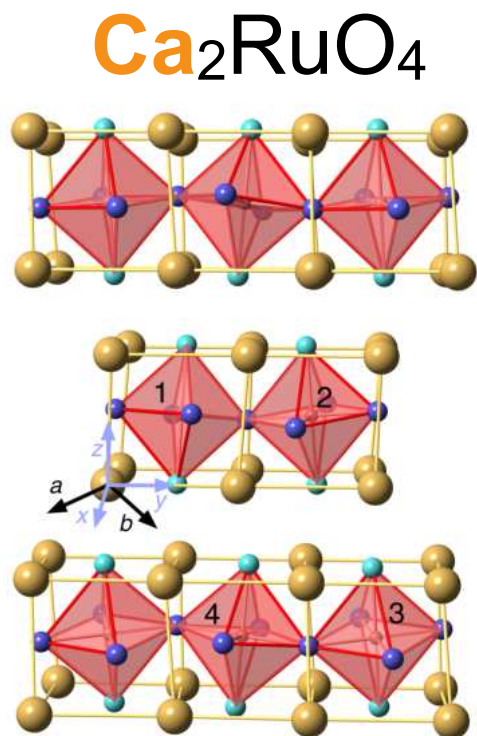
the Sr_2RuO_4 family: exploit complexity



some known facts..

replacing Sr with Ca: metal-insulator transition

$Ru\ 4d\ t_{2g}^4$



Nature of Mott transition

PRL **104**, 226401 (2010)

PHYSICAL REVIEW LETTERS

week ending
4 JUNE 2010

Nature of the Mott Transition in Ca_2RuO_4

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²*I. Institut für Theoretische Physik, Universität Hamburg, Jungiusstraße 9, D-20355 Hamburg, Germany*

(Received 12 January 2010; published 1 June 2010)

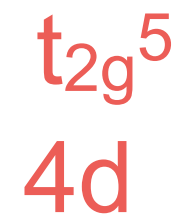
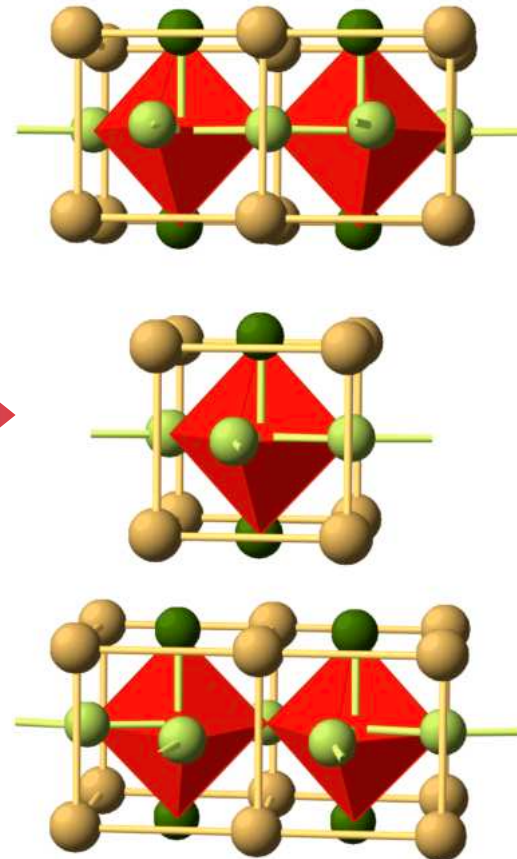
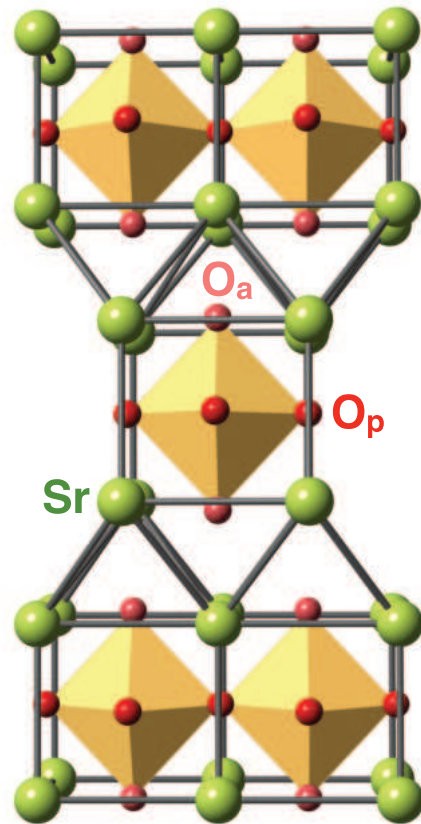
We study the origin of the temperature-induced Mott transition in Ca_2RuO_4 . As a method we use the local-density approximation + dynamical mean-field theory. We show the following. (i) The Mott transition is driven by the change in structure from long to short **c**-axis layered perovskite ($L\text{-}Pbca \rightarrow S\text{-}Pbca$); it occurs together with orbital order, which follows, rather than produces, the structural transition. (ii) In the metallic $L\text{-}Pbca$ phase the orbital polarization is ~ 0 . (iii) In the insulating $S\text{-}Pbca$ phase the lower energy orbital, $\sim xy$, is full. (iv) The spin-flip and pair-hopping Coulomb terms reduce the effective masses in the metallic phase. Our results indicate that a similar scenario applies to $\text{Ca}_{2-x}\text{Sr}_x\text{RuO}_4$ ($x \leq 0.2$). In the metallic $x \leq 0.5$ structures electrons are progressively transferred to the xz/yz bands with increasing x ; however, we find no orbital-selective Mott transition down to ~ 300 K.

$$\epsilon_{CF} + \Delta \epsilon_{CF} \quad xy^2xz^1yz^1$$

λ (spin-orbit coupling) **not** essential for MIT, but key to magnetism

Phys. Rev. B **95**, 075145 (2017); Phys. Rev. B **101**, 205128 (2020)

replacing Ru with Rh



rotations: xy-band FS missing

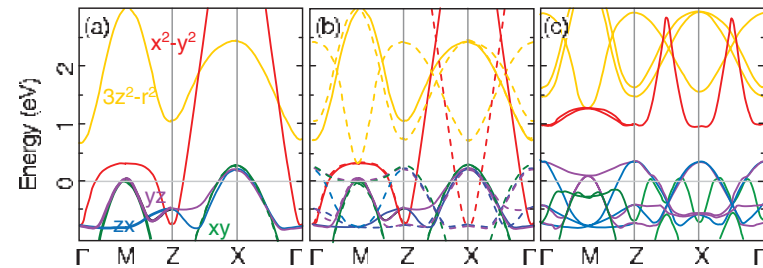
PRL **97**, 106401 (2006)

PHYSICAL REVIEW LETTERS

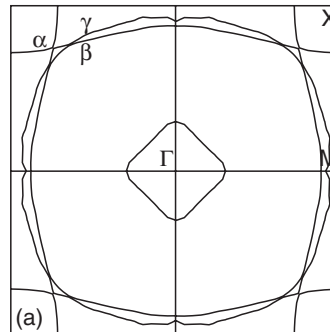
week ending
8 SEPTEMBER 2006

Missing *xy*-Band Fermi Surface in 4*d* Transition-Metal Oxide Sr_2RhO_4 : Effect of the Octahedra Rotation on the Electronic Structure

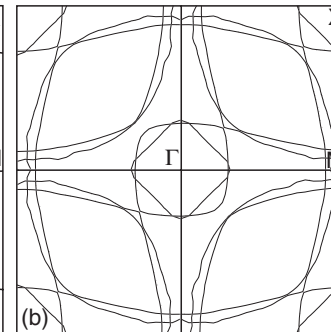
B. J. Kim,¹ Jaejun Yu,¹ H. Koh,² I. Nagai,³ S. I. Ikeda,³ S.-J. Oh,¹ and C. Kim⁴



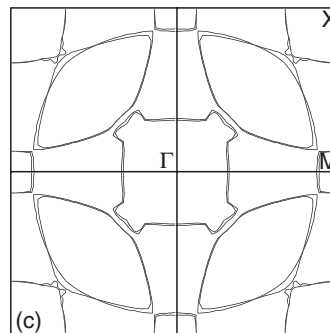
undistorted



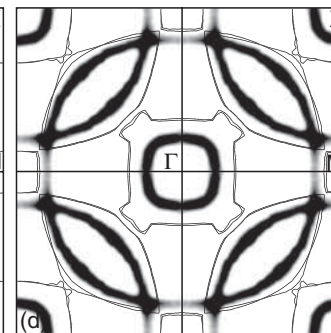
undistorted+folded



distorted



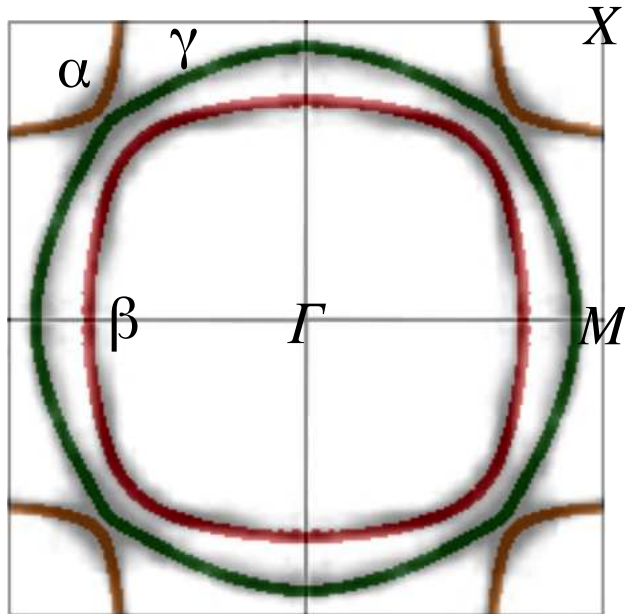
distorted
vs experiments



Ru $\rightarrow$ Rh: moderate masses, xy band full

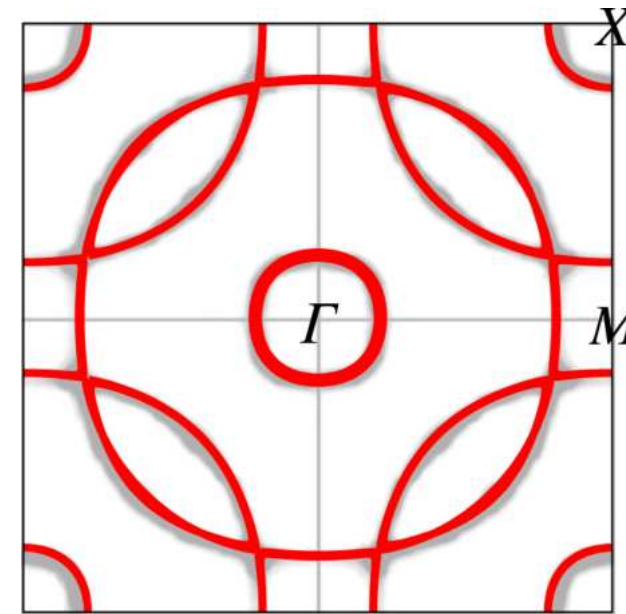
t_{2g}^4

Sr_2RuO_4 :



t_{2g}^5

Sr_2RhO_4 :



$\lambda + \Delta\lambda$

theory:

G. Zhang, E. Gorelov, E. Sarvestani, and E. Pavarini, Phys. Rev. Lett. **116**, 106402 (2016)

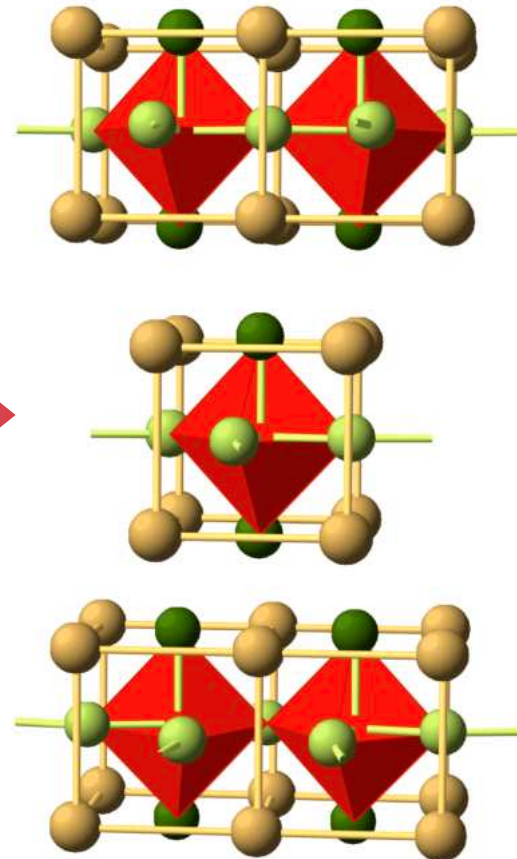
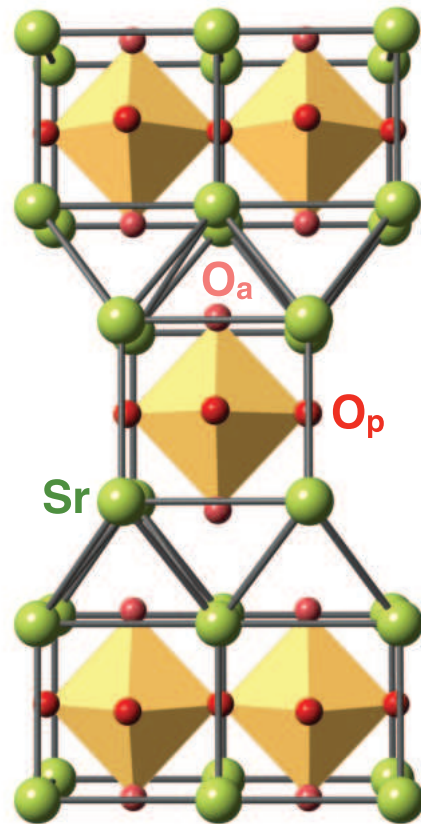
G. Zhang and E. Pavarini, Phys. Rev. B **99**, 125102 (2019)

experiments:

A. Damascelli, D. H. Lu, K. M. Shen, N. P. Armitage, F. Ronning, D. L. Feng, C. Kim, Z.-X. Shen, T. Kimura, Y. Tokura, Z. Q. Mao, and Y. Maeno, Phys. Rev. Lett. **85**, 5194 (2000)

B. J. Kim, J. Yu, H. Koh, I. Nagai, S. I. Ikeda, S.-J. Oh, and C. Kim, Phys. Rev. Lett. **97**, 106401 (2006)

replacing Ru with Ir



Sr₂IrO₄: t_{2g} analogous of HTSCs?

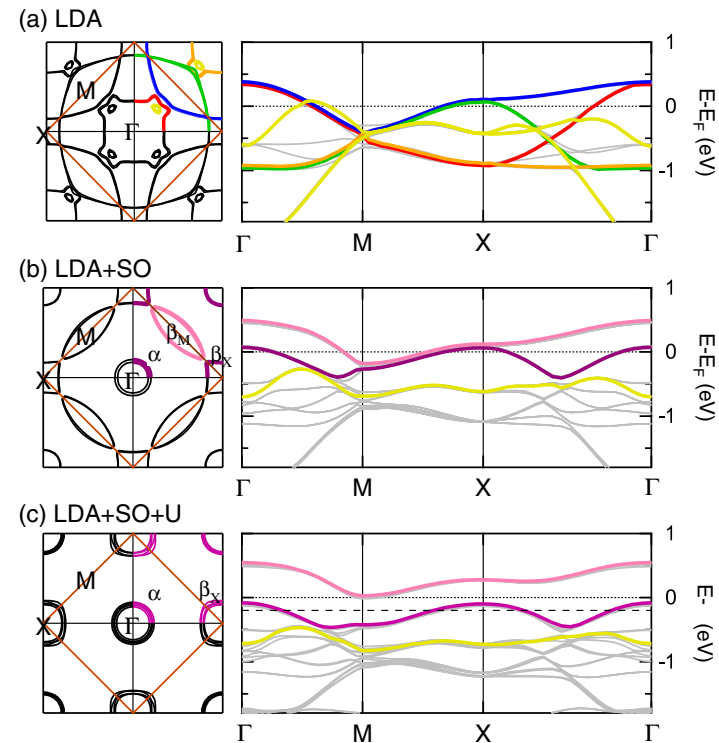
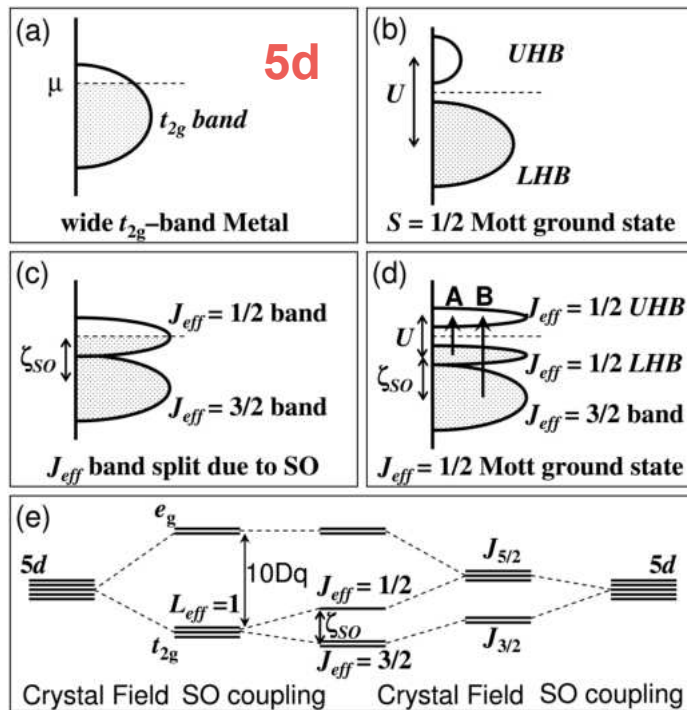
PRL **101**, 076402 (2008)

PHYSICAL REVIEW LETTERS

week ending
15 AUGUST 2008

Novel $J_{\text{eff}} = 1/2$ Mott State Induced by Relativistic Spin-Orbit Coupling in Sr₂IrO₄

B. J. Kim,¹ Hosub Jin,¹ S. J. Moon,² J.-Y. Kim,³ B.-G. Park,⁴ C. S. Leem,⁵ Jaejun Yu,¹ T. W. Noh,² C. Kim,⁵ S.-J. Oh,¹
J.-H. Park,^{3,4,*} V. Durairaj,⁶ G. Cao,⁶ and E. Rotenberg⁷

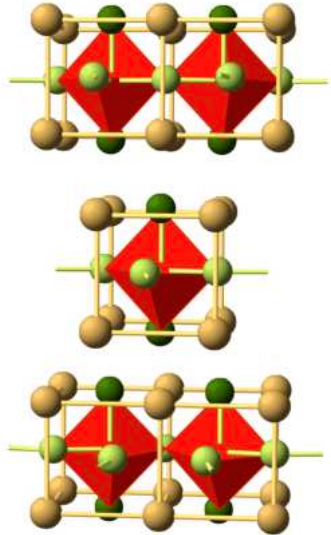


angle-resolved photoemission, optical conductivity,

x-ray absorption measurements

no, rather $j_{\text{eff}}=1/2$ spin-orbital order

spin-orbit large but non sufficiently for 1-band picture



PHYSICAL REVIEW LETTERS **131**, 036504 (2023)

Multiorbital Nature of Doped Sr_2IrO_4

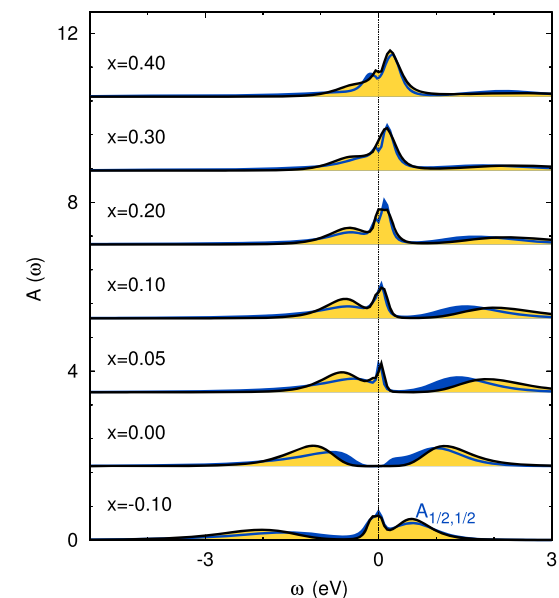
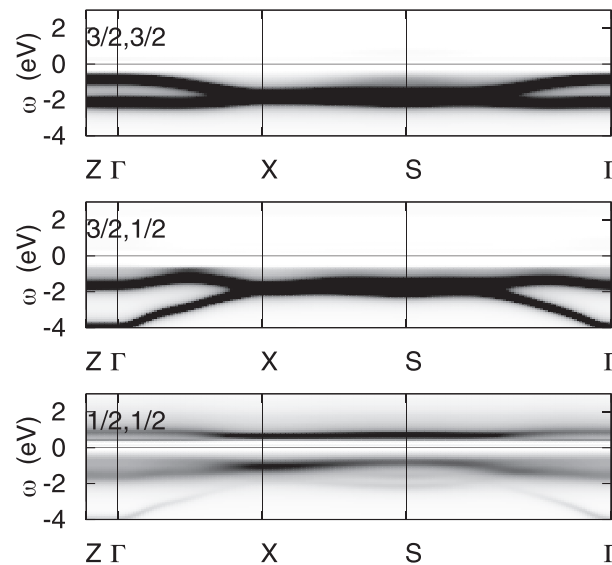
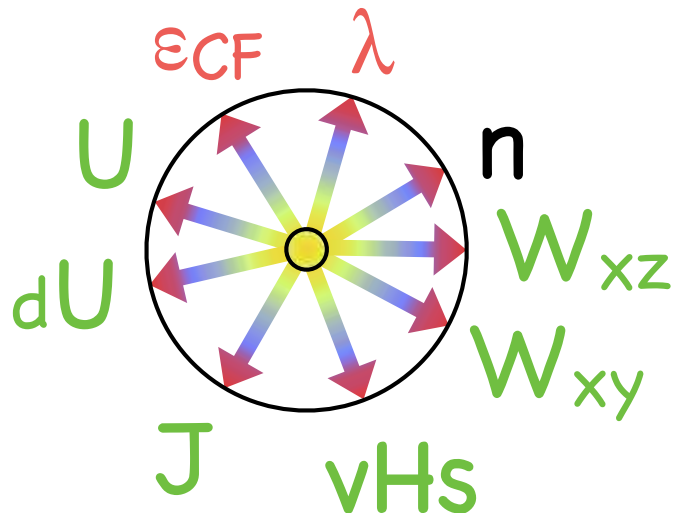
Guoren Zhang^{1,2} and Eva Pavarini^{3,4}

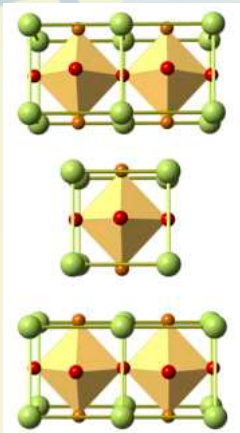
¹School of Sciences, Nantong University, Nantong, 226019, People's Republic of China

²Key Laboratory of Materials Physics, Institute of Solid State Physics, HFIPS, Chinese Academy of Sciences, Hefei 230031, People's Republic of China

³Institute for Advanced Simulation, Forschungszentrum Jülich, 52425 Jülich, Germany

⁴JARA High-Performance Computing, Forschungszentrum Jülich, 52425 Jülich, Germany



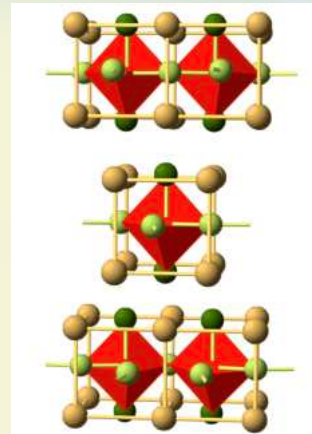


correlated metal

$$m_{xy} > m_{xz/yz}$$

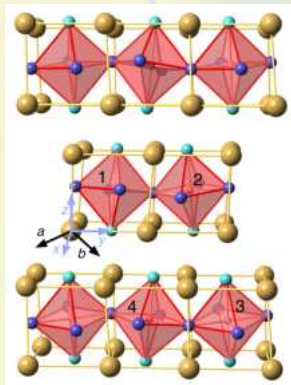


$$(t_{2g})^n$$



weakly correlated metal

only 2-bands
xz/yz at Fermi surface

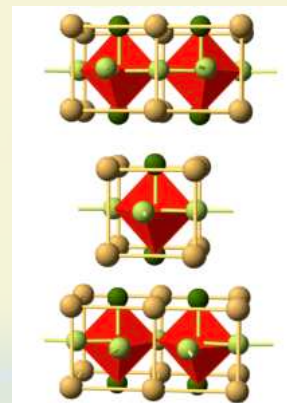


small gap insulator

xy-orbital order

spin-orbit important
only for magnetism

$$(t_{2g})^4$$



spin-orbit Mott insulator
small gap

$j=1/2$ orbital order
only 1-band half-filled

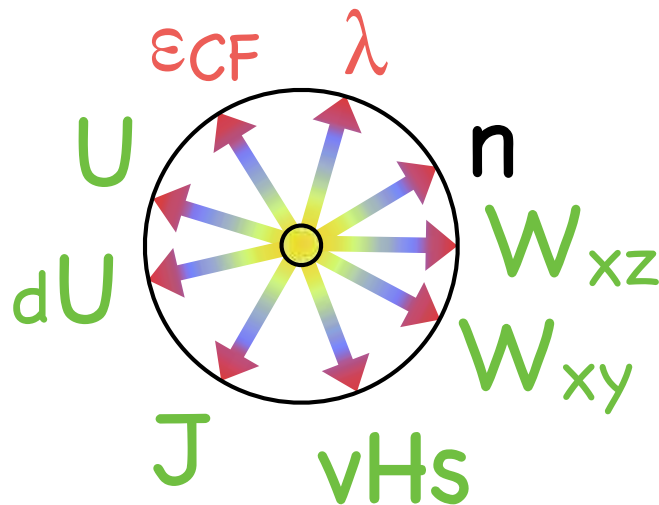
$$(t_{2g})^5$$

can we build a unifying picture?

yes!

ϵ_{CF} changes natural occupations

build map of CF (ϵ_{CF}) effects!



PHYSICAL REVIEW LETTERS **132**, 236505 (2024)

Map of Crystal-Field Effects in Correlated Layered t_{2g}^n Perovskites

Neda Samani,¹ Guoren Zhang^{1,2,3} and Eva Pavarini^{1,4}

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²School of Physics and Technology, *Nantong University*, Nantong 226019, People's Republic of China

³Key Laboratory of Materials Physics, Institute of Solid State Physics, *HFIPS*, Chinese Academy of Sciences, Hefei 230031, People's Republic of China

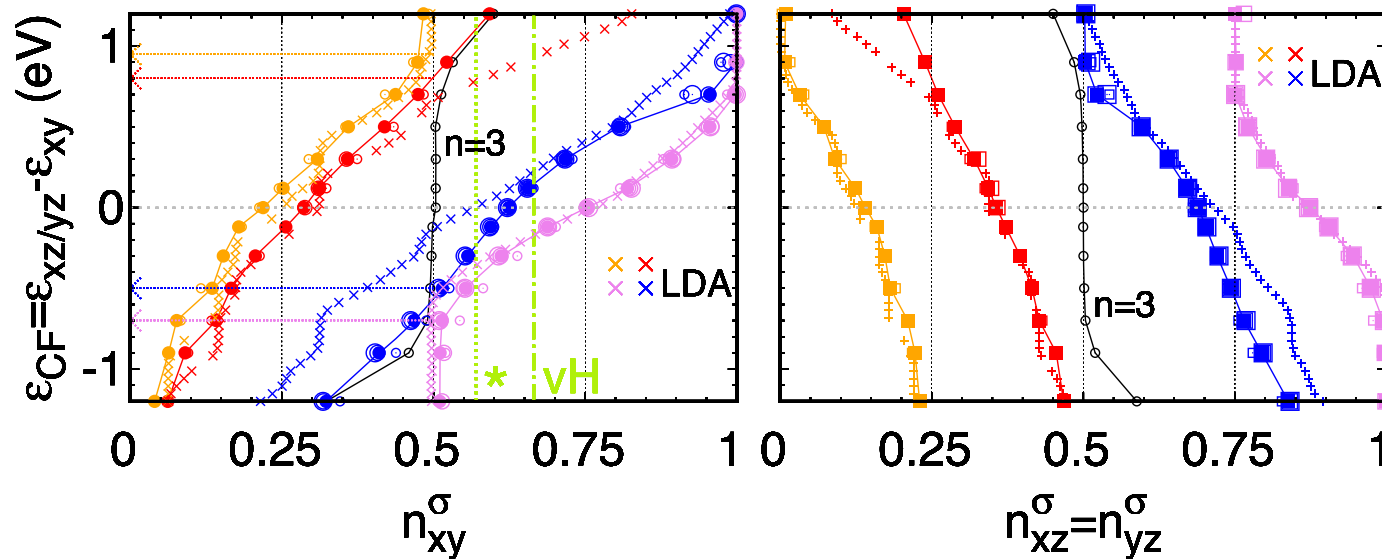
⁴JARA High-Performance Computing, *Forschungszentrum Jülich*, Germany

(natural orbitals=eigenvalues of occupation matrix)

ϵ_{CF} changes natural occupations

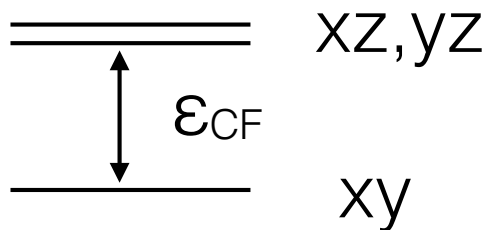
$\lambda = 0$ natural orbitals= $xy, xz/yz$

t_{2g}^n



$n=1$
 $n=2$
 $n=4$
 $n=5$

CF splitting



crosses: $U=0$

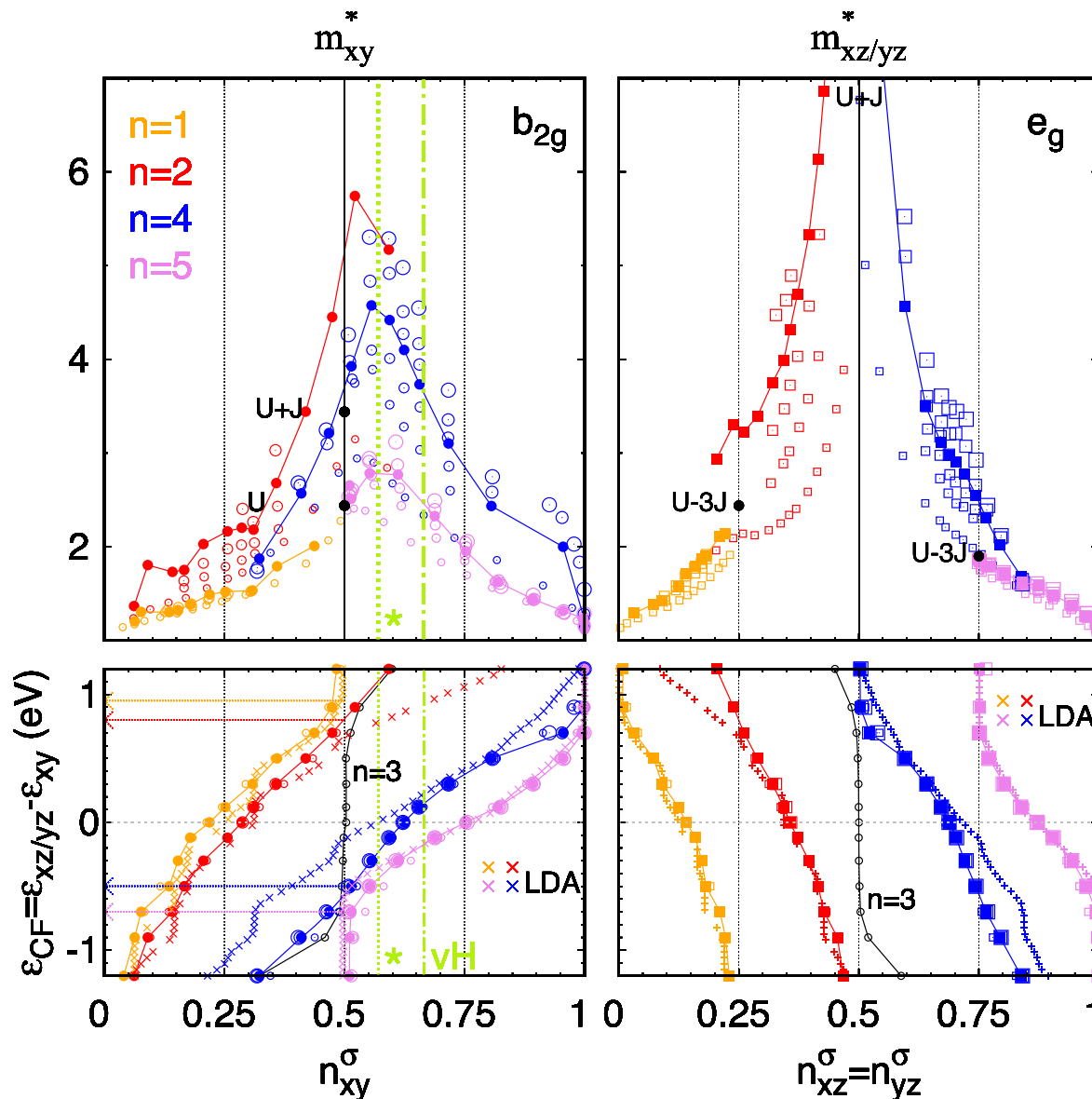
closed symbols: finite U

$$n = n_{xy} + 2n_{xz/yz}$$

map of crystal-field splitting: effective masses

$$t_{2g}^n$$

$$\lambda = 0$$



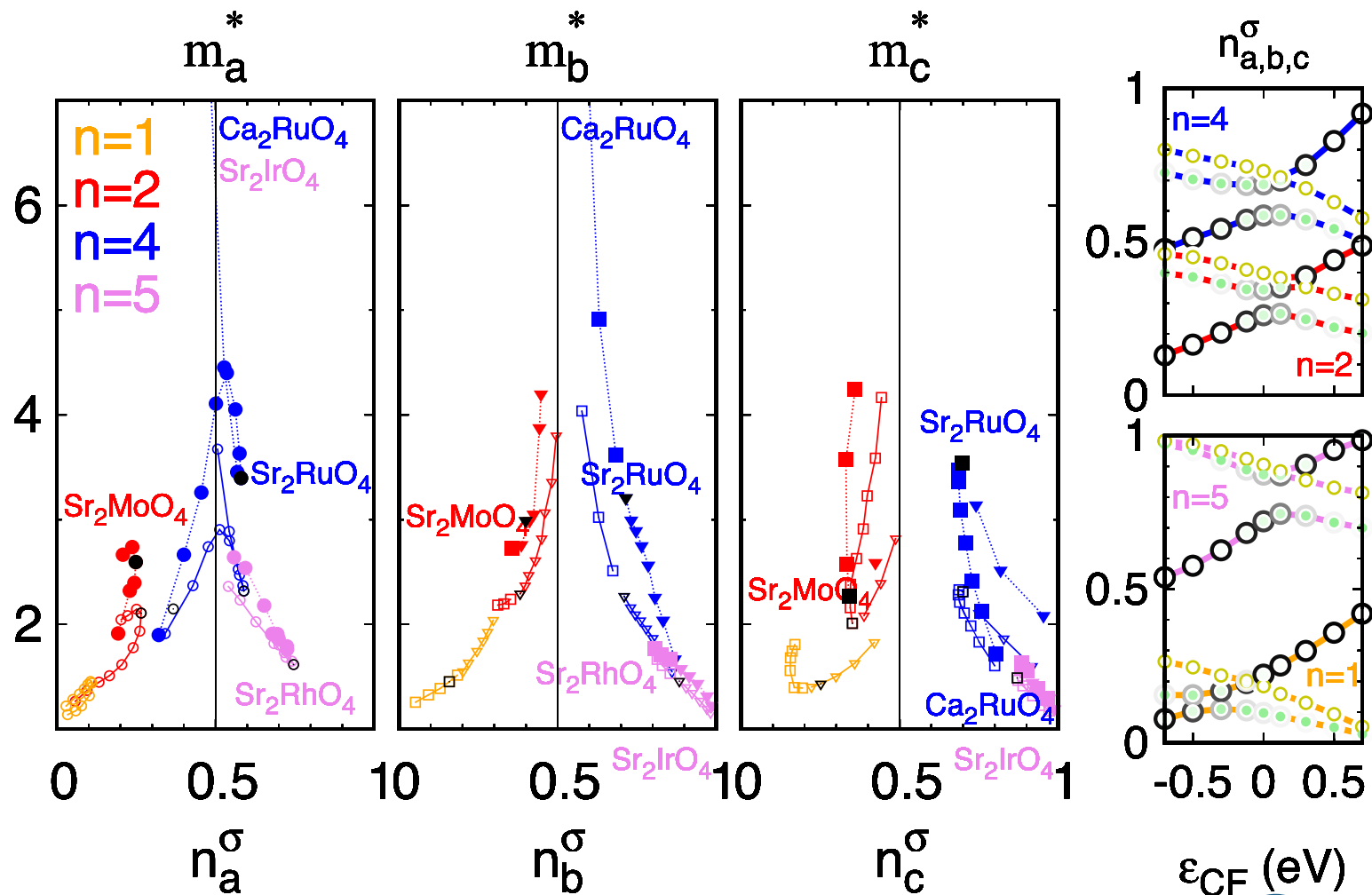
m^* diverges:
insulator

$m^*=1$: not
correlated

moderate (realistic) U, J , fixed U/W

switch on moderate spin-orbit coupling

λ as in Sr_2RuO_4 (4d)

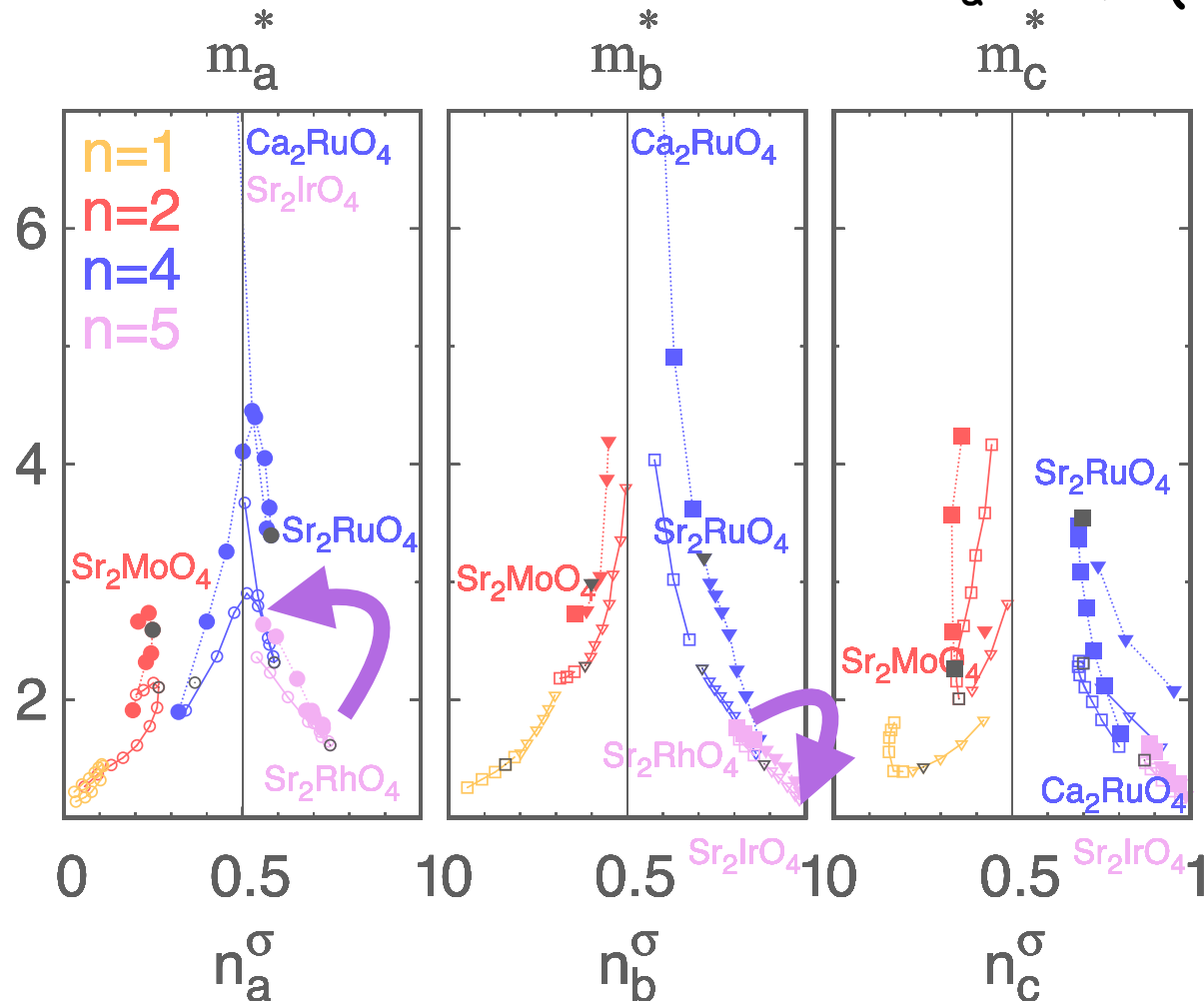


putting the map to a test: Sr_2IrO_4

λ as in Sr_2RuO_4 (4d)

change of configuration effect:

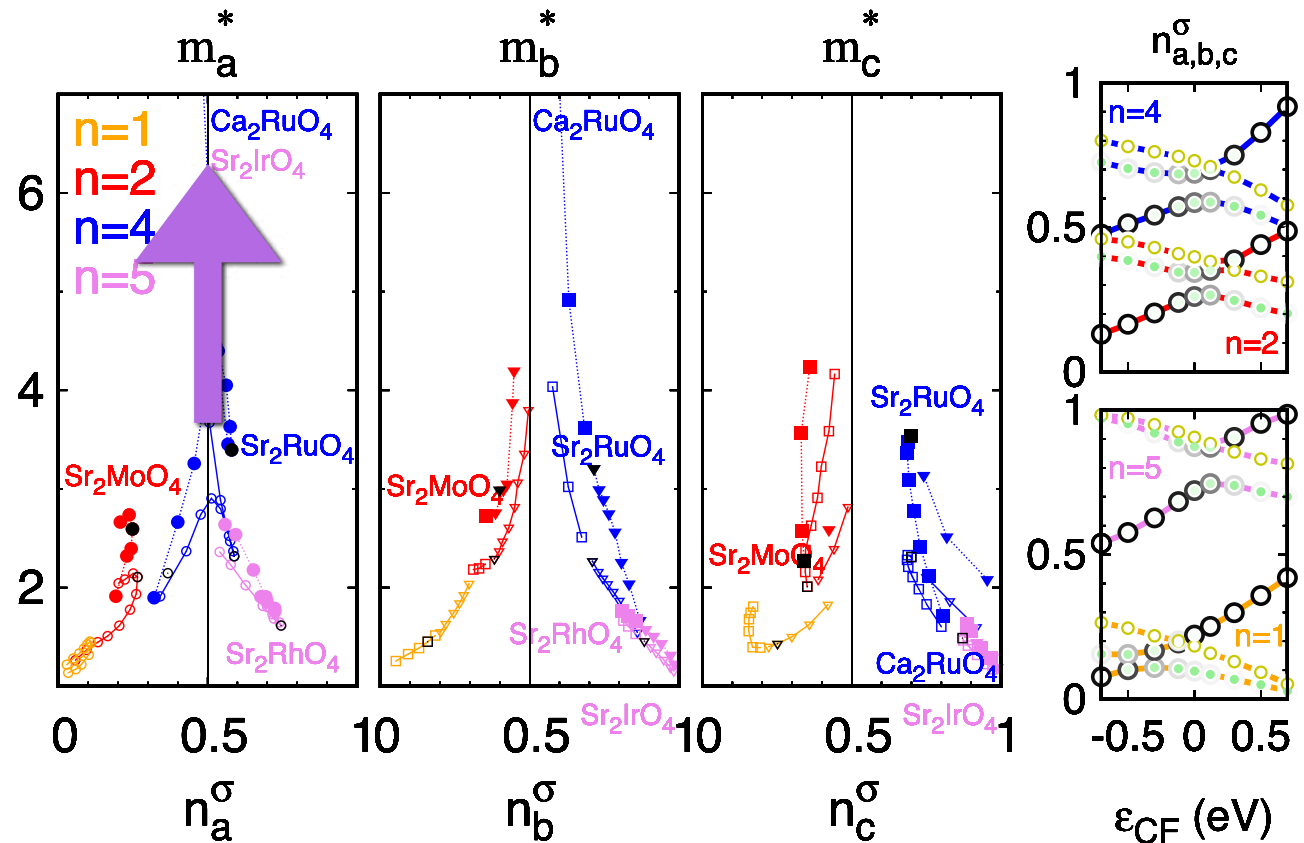
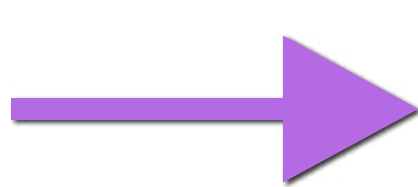
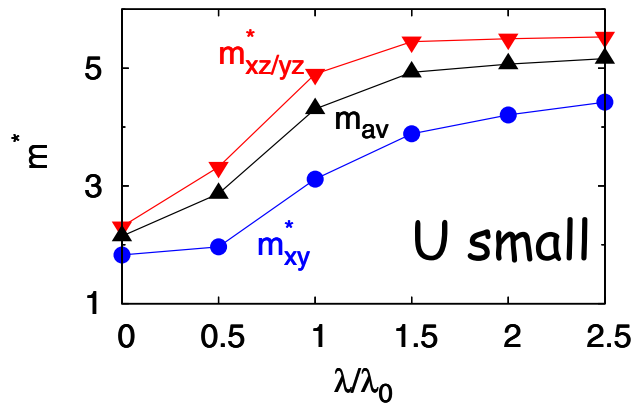
mass-matrix trace invariant: $m_a^* \sim 1 + 3(m_{av} - 1)$



putting the map to a test: Sr_2IrO_4

λ two times larger (5d)

suppression of spin-orbital fluctuations



putting the map to a test: pressure effects

Sr_2IrO_4 : persistence of insulating state
state under very high pressure
(tetragonal till 40 GPa)

IOP Publishing

Journal of Physics: Condensed Matter

J. Phys.: Condens. Matter **26** (2014) 255603 (5pp)

doi:10.1088/0953-8984/26/25/255603

Persistent non-metallic behavior in Sr_2IrO_4 and $\text{Sr}_3\text{Ir}_2\text{O}_7$ at high pressures

D A Zocco^{1,6}, J J Hamlin^{1,7}, B D White¹, B J Kim^{2,3}, J R Jeffries⁴, S T Weir⁴,
Y K Vohra⁵, J W Allen³, and M B Maple¹

PHYSICAL REVIEW B **101**, 144102 (2020)

Persistent insulating state at megabar pressures in strongly spin-orbit coupled Sr_2IrO_4

Chunhua Chen,^{1,2} Yonghui Zhou,^{1,*} Xuliang Chen,¹ Tao Han,³ Chao An,³ Ying Zhou,³ Yifang Yuan,^{1,2} Bowen Zhang,^{1,2}
Shuyang Wang,^{1,2} Ranran Zhang,¹ Lili Zhang,⁴ Changjin Zhang,^{1,3,5} Zhaorong Yang,^{1,3,5,†}
Lance E. DeLong,⁶ and Gang Cao^{7,‡}

¹Anhui Province Key Laboratory of Condensed Matter Physics at Extreme Conditions, High Magnetic Field Laboratory,
Chinese Academy of Sciences, Hefei 230031, China

²Science Island Branch of Graduate School, University of Science and Technology of China, Hefei 230026, China

³Institutes of Physical Science and Information Technology, Anhui University, Hefei 230601, China

⁴Shanghai Synchrotron Radiation Facility, Shanghai Advanced Research Institute, Chinese Academy of Sciences, Shanghai 201204, China

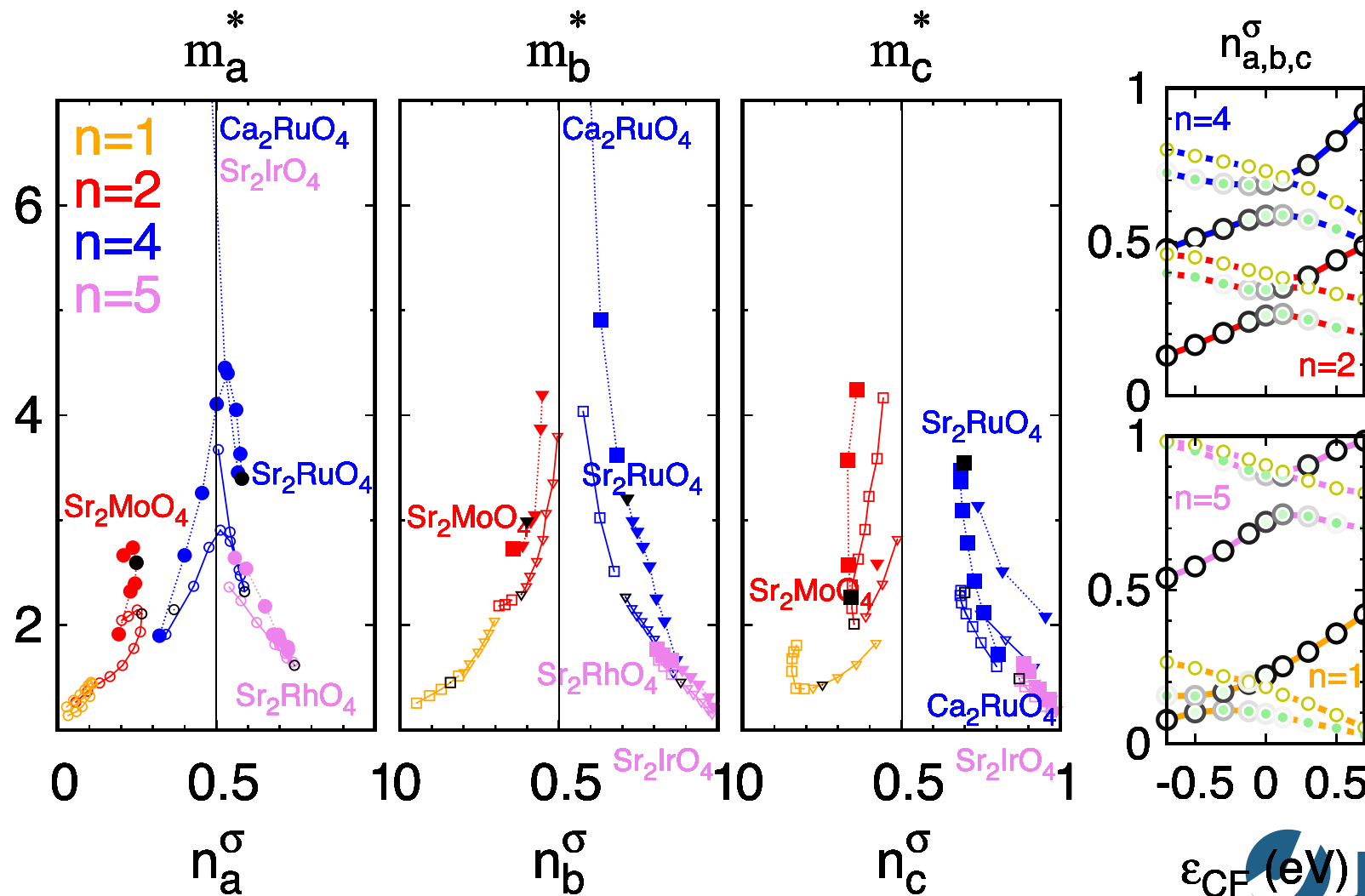
⁵Collaborative Innovation Center of Advanced Microstructures, Nanjing University, Nanjing 210093, China

⁶Department of Physics and Astronomy, University of Kentucky, Lexington, Kentucky 40506, USA

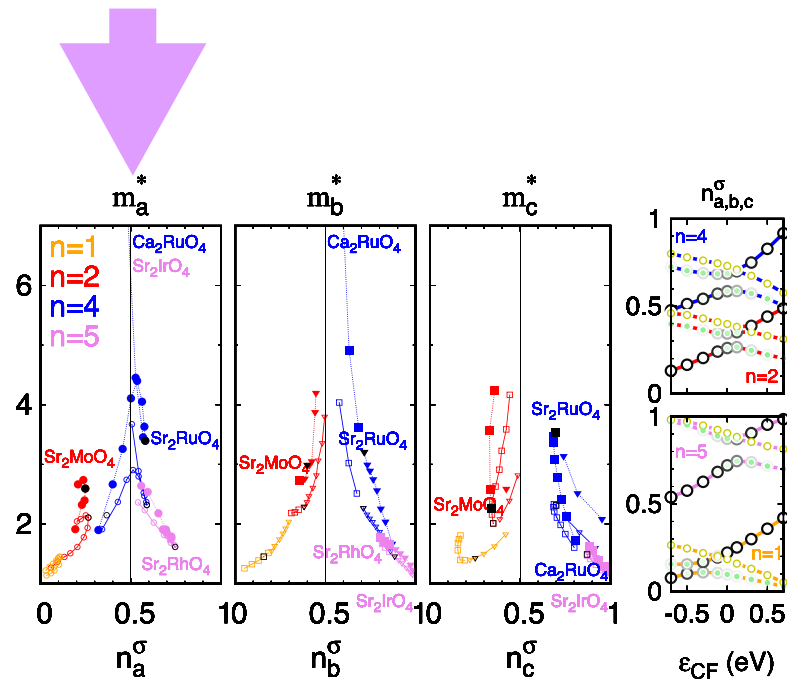
⁷Department of Physics, University of Colorado at Boulder, Boulder, Colorado 80309, USA

position in the map

$j_{\text{eff}}=1/2$ orbital order is stable

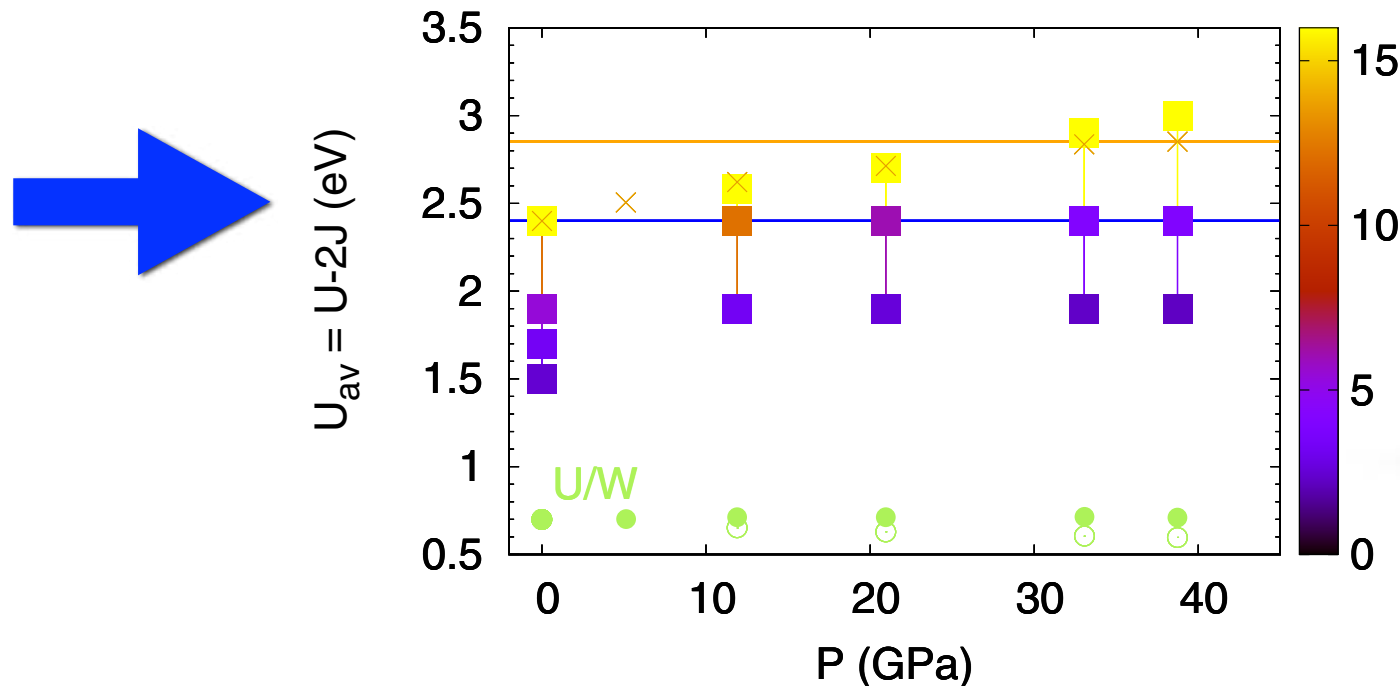


W increases but slowed down by rotations



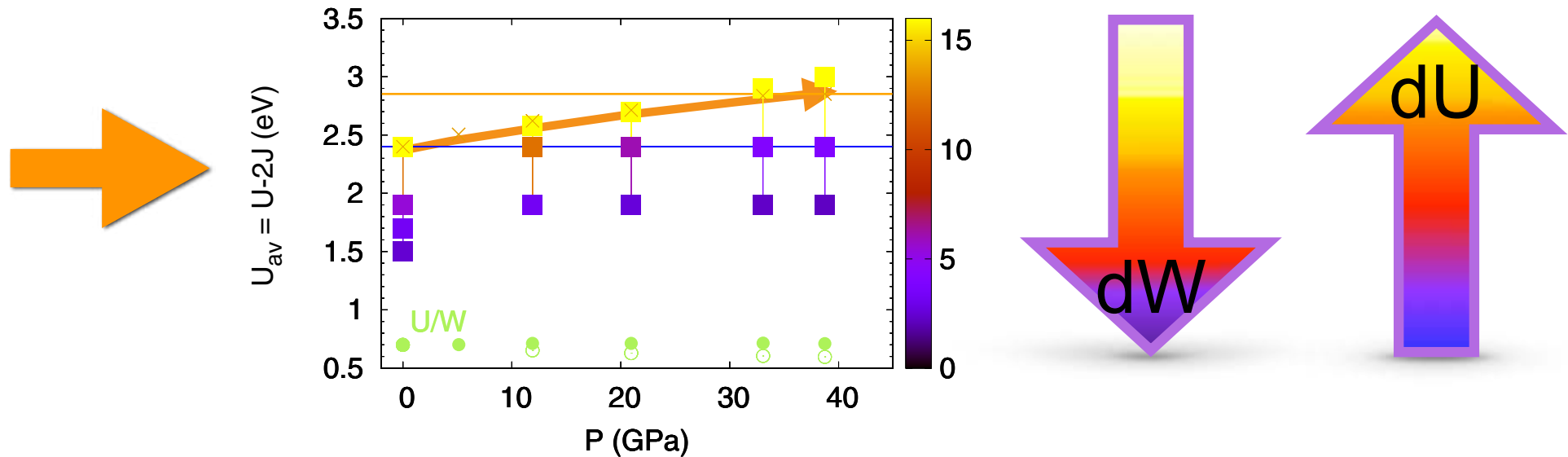
metal already at 12 GPa

color map:
effective mass



however $W/U \sim \text{constant}$

screened U (cRPA) also increases



PHYSICAL REVIEW B **111**, L121101 (2025)

Letter

Nature of the high-pressure insulating state in Sr_2IrO_4 : Mott picture

Guoren Zhang^{1,2,3}, Hanif Hadipour⁴, and Eva Pavarini³

¹School of Physics and Technology, [Nantong University](#), Nantong 226019, People's Republic of China

²Key Laboratory of Materials Physics, [Institute of Solid State Physics, HFIPS](#),
Chinese Academy of Sciences, Hefei 230031, People's Republic of China

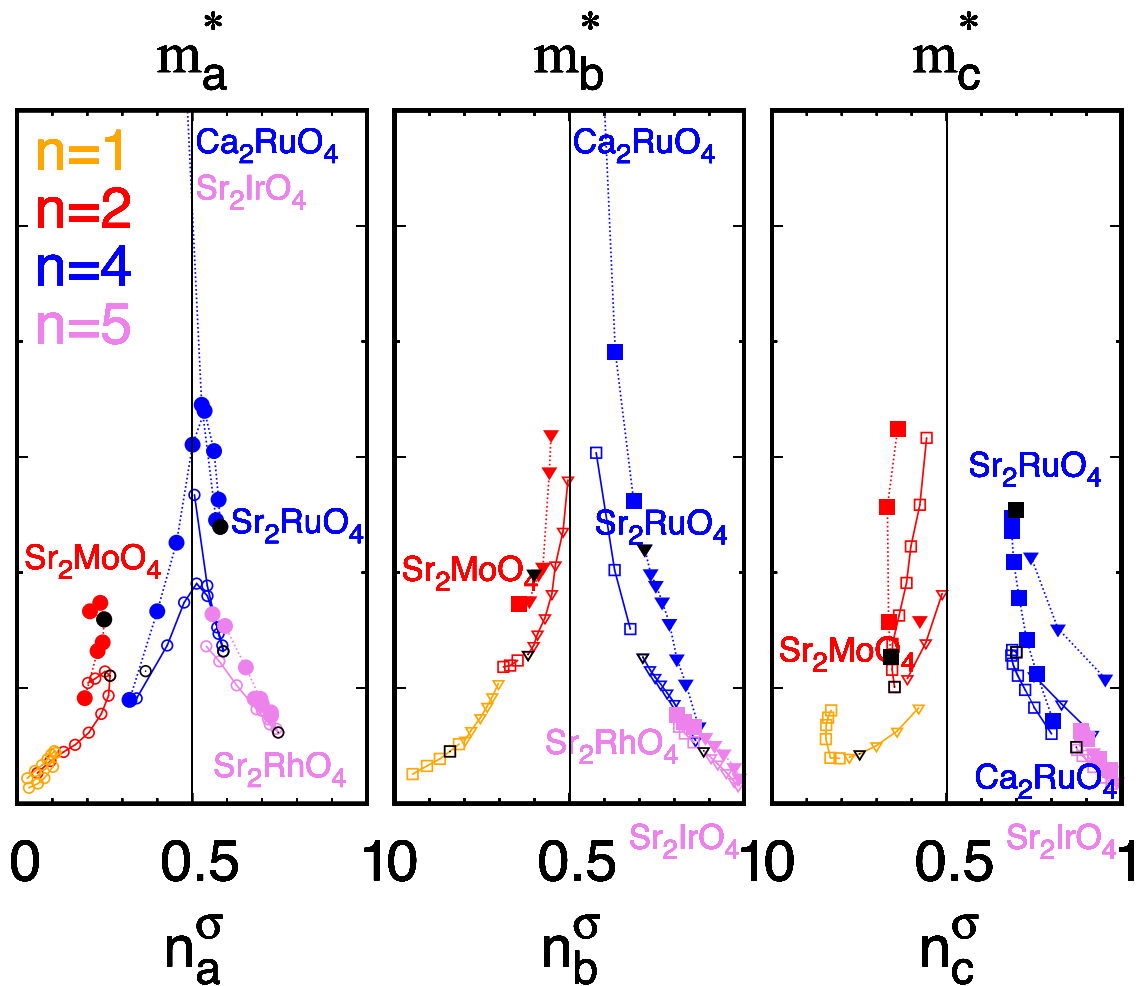
³Peter Grünberg Institute, [Forschungszentrum Jülich](#), D-52425 Jülich, Germany

⁴Department of Physics, [University of Guilan](#), Rasht 41335-1914, Iran

very close to MIT under high pressure!

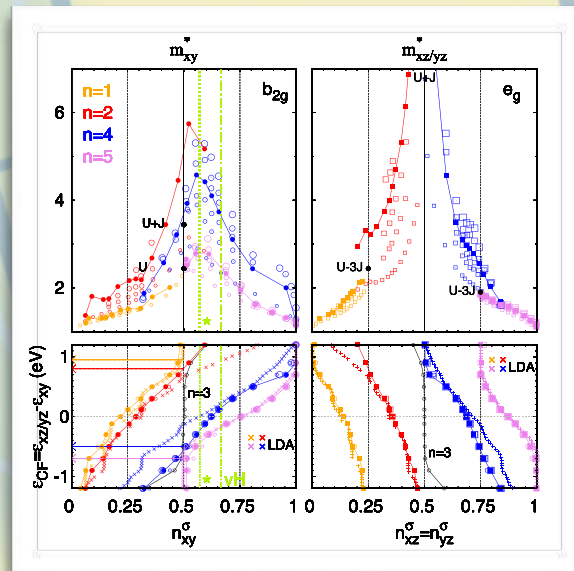
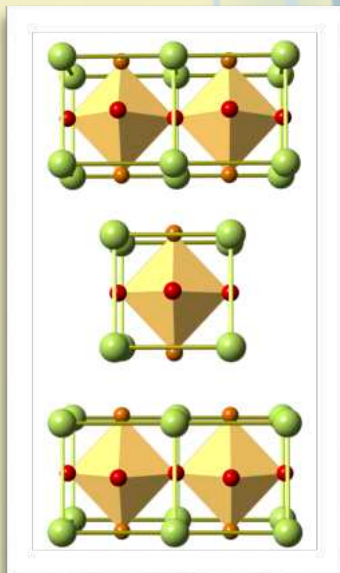
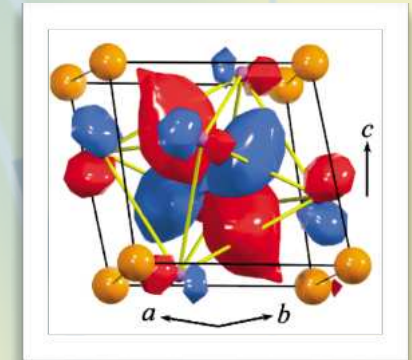
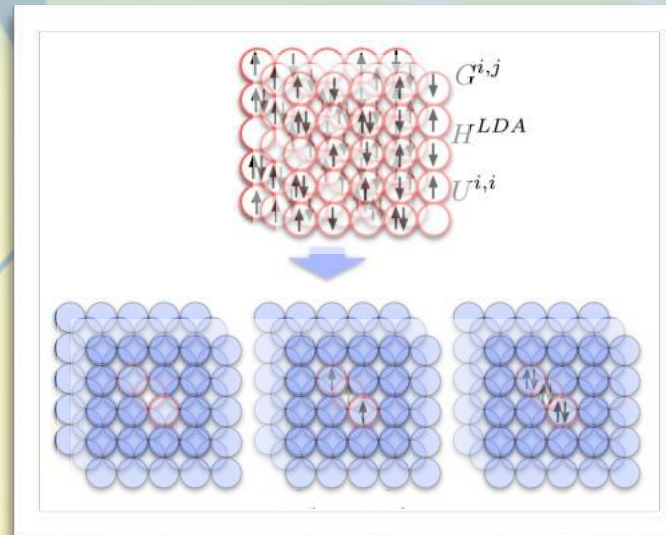


combined effect of stability of $j=1/2$
orbital-order and constant W/U



conclusions

- correlated materials
- DFT+DMFT



- Sr_2RuO_4 and its family
- disentangle complexity via an unifying map

main people

E. Gorelov

E. Sarvestani

H. Hadipour

G. Zhang

N. Samani

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Jureca-DC & Juwels

Gauss Centre for Supercomputing
and VSR-JARA at FZJ

thank you!